

ATIPROTON - A TOOL FOR NUCLEAR STUDIES

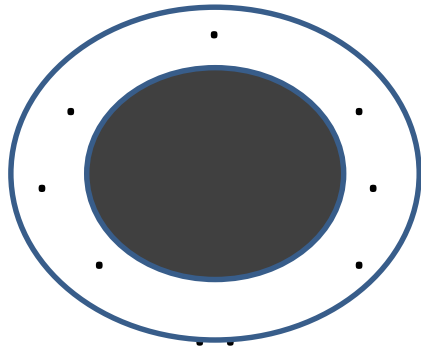
Sławomir Wycech

SPECIAL FEATURES OF ANTI-PROTON

STRONG ABSORPTION IN NUCLEI

$$\sigma_{\text{abs}} \sim 200 \text{ mb}$$

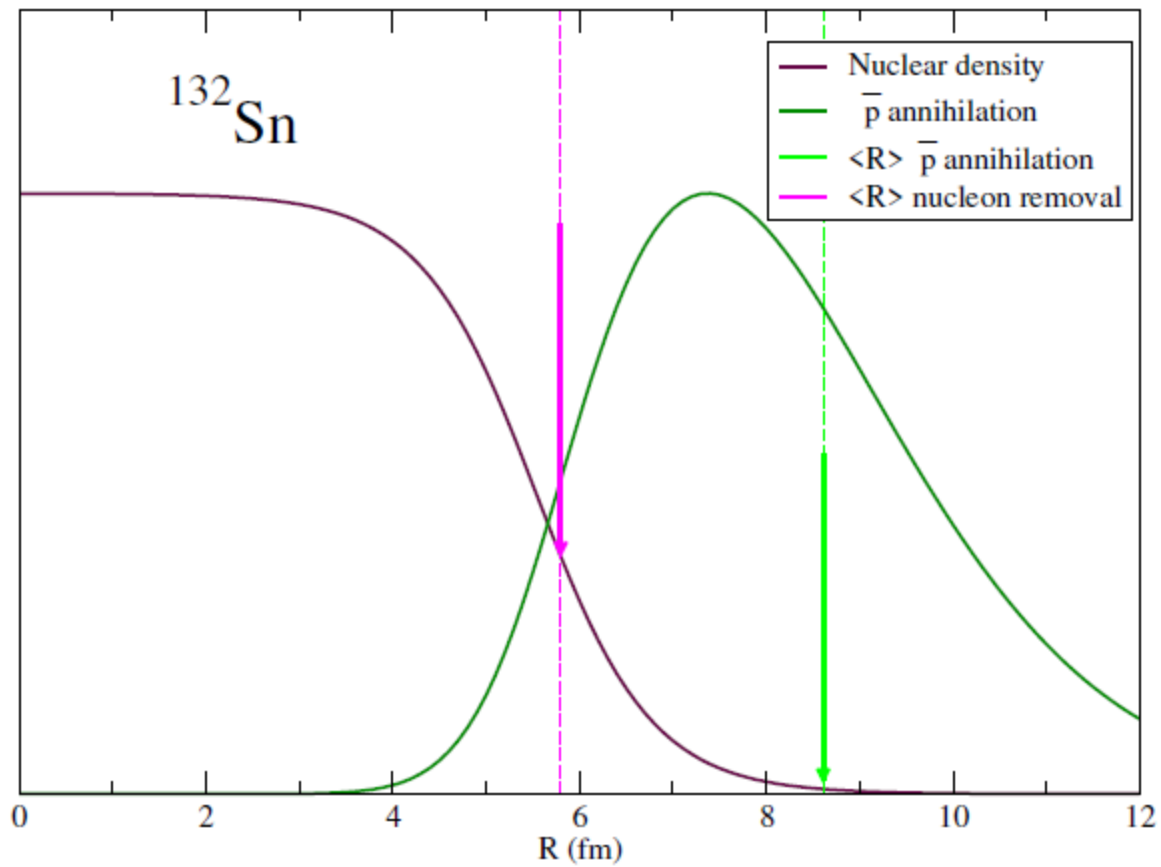
$$\text{free path in nuclei } 1/\sigma\rho = 0.3 \text{ fm}$$



does not enter nuclei

CHARACTERISTIC TRACE $N \bar{N} \rightarrow \pi.. \pi.. \pi.. \pi$ 2 – 8 mesons
2/3 charged

$Z = 50$, $N = 88$: fancy, unstable nucleus to study by PUMA
Expected atomic – nuclear density overlap



WHY NUCLEAR SURFACE IS INTERESTING

* Symmetry energy

$$\beta = (N - Z)/A,$$

$$\frac{E}{A}(\rho, \beta) = \frac{E}{A}(\rho, 0) + S_N(\rho)\beta^2 + \dots \quad \text{n,p Fermi Gas} \quad S_N = \frac{1}{3}E_F$$

ρ = density

Droplet Model

$$E_{(\text{binding})} / A = a_v \quad - \quad S_N \beta^2 \quad + \quad \dots$$

attractive repulsive due to Pauli

THESE CANCEL AT NUCLEAR SURFACE WITH THE INCREASING NEUTRON/PROTON RATIO ? NUCLEAR MODEL DEPENDENT

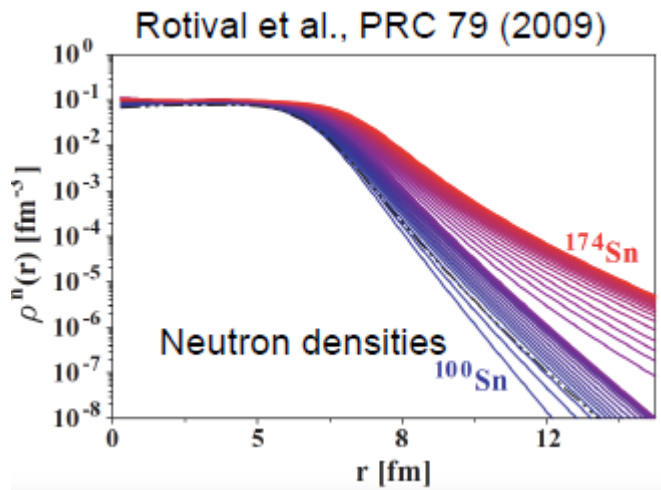
** ARE THERE (np.) or (nnpp) CORRELATIONS AT DISTANT SURFACE.

*** WHAT IS THE FERMİ MOMENTUM AT SURFACE

EXPECTATION - an EXAMPLE

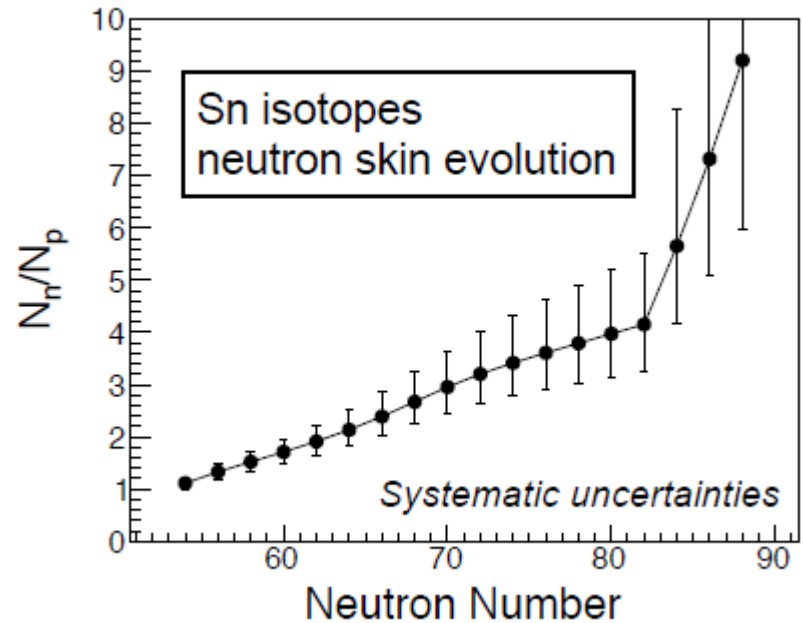
from A. Obertelli

Neutron tails



Pbar nuclear absorption region

n/p ratio expected at capture radius



MOTIVATION

PUMA at CERN : From
Alexandre Obertelli



Fig. 7: Itinerary of PUMA from ELENA to ISOLDE.

PUMA EXPERIMENT

PRODUCES ANTIPROTONS

TRANSPORTS in A „BOTTLE”

COLLIDES ANTIPROTONS WITH UNSTABLE NUCLEI

MAKES ANTIPROTONIC ATOMS

waits for X ray cascade, and nuclear capture

DETECTS CHARGED π MESONS FROM ANNIHILATION

ESTABLISH ATOMIC ORBITS OF NUCLEAR CAPTURE

done

CALCULATE ABSORPTION RATIO

in progress

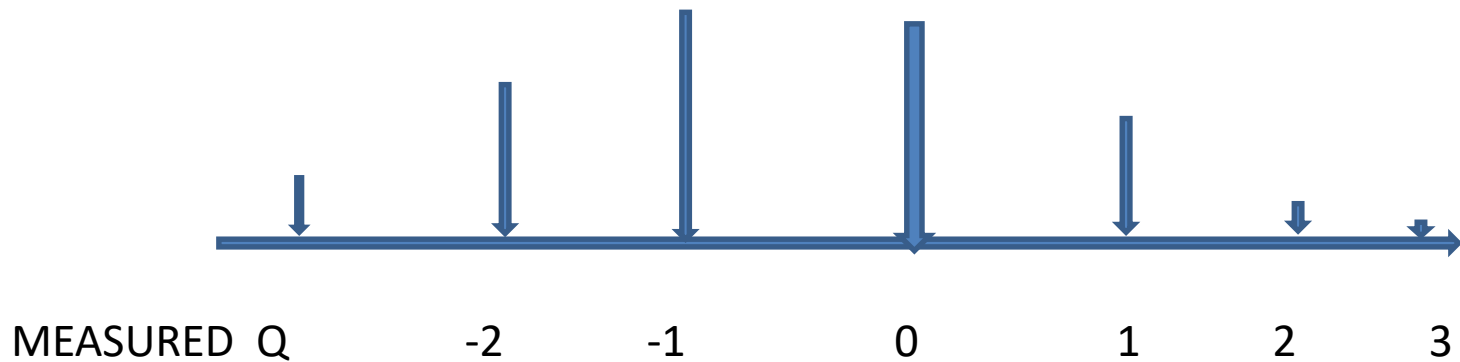
$$\sigma (\text{Pbar} - n) / \sigma(\text{Pbar} - p)$$

→ NEUTRON HALO (SKIN)

→ NN , PPNN CORRELATIONS ON SURFACE

EXTRACTION OF CAPTURE ORBITALS FROM TOTAL MESONIC CHARGE

INITIAL $Q = 0$ capture on proton 5 mesons mitted
 $Q = -1$ capture on neutron 5 mesons emitted



ANALYSIS OF FINAL STATE MESONIC REACTIONS

(0) CHOSE PARAMETERS FOR ABSORPTION $\pi NN \rightarrow NN$,
CHARGE EXCHANGE $\pi^+ \rightarrow \pi^0$, $\pi^0 \rightarrow \pi^+$
 $\pi^- \rightarrow \pi^0$, $\pi^0 \rightarrow \pi^-$

(1) FIT PARAMETERS TO P(Q) DATA

(2) CALCULATE PARAMETERS

(3) COMPARE FITTED TO CALCULATED



extract the orbits of captures



calculate neutron haloes

CAPTURE ORBITS
 IN PIONISATION
 MEASUREMENTS
 VERSUS X RAY DATA

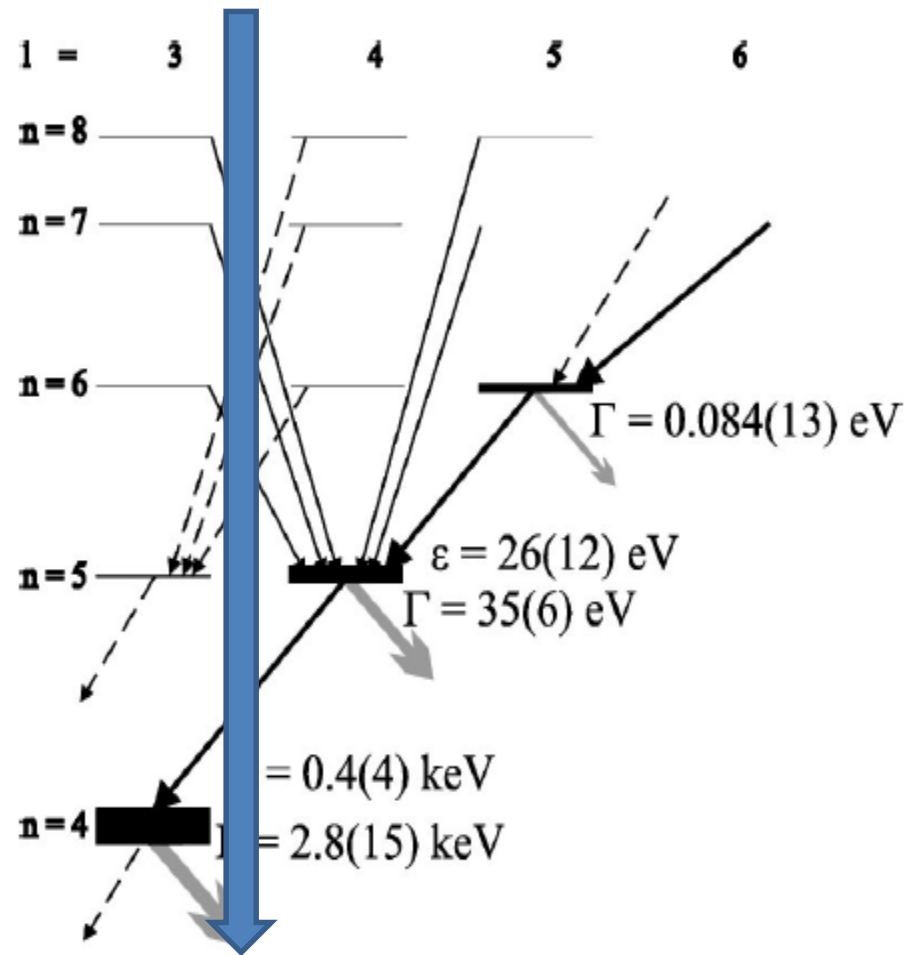


FIG. 3. Mean widths and shifts of all levels with measurable strong interaction effects. The weight of the different calcium iso-

RESULTS

OLD DATA : N, C, Ti, Ta, Pb analysed

S.W.,K.P. Phys Rev. C (2023) 108

DOMINANT CAPTURE ORBITS :
THE LOWEST STATES REACHED IN ATOMIC CASCADE

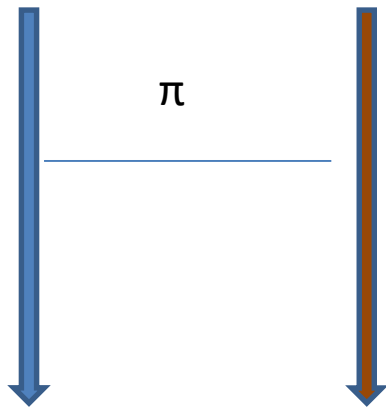
Rms RADII OF NEUTRON DENSITIES CONSISTENT
WITH OTHER EXPERIMENTS

SIZABLE ERRORS -

uncertain ratio of captures on protons relative to neutrons

SECOND ESSENTIAL JOB FOR THEORY

A MODEL FOR NUCLEON-ANTINUCLEON INTERACTIONS



- + π - π CORRELATED BY DISPERSION RELATIONS
- + PHENOMENOLOGY MODEL DEPENDENT

EXTEND BEST MODEL TO LOW ENERGY,
FIND BOUND STATES

COMPARISON N-Nbar Interaction OF MODELS

J.Carbonell, G.Hupin. S.W. : EPJA59(2023)259

IMPROVING N- Nbar INTERACTION POTENTIAL

DATA CROSS SECTIONS ONLY , MANY PARTIAL WAVES (NO PAULI)
NO LOW ENERGY DATA

WHAT HAPPENS
BELOW 200 MV/C ?

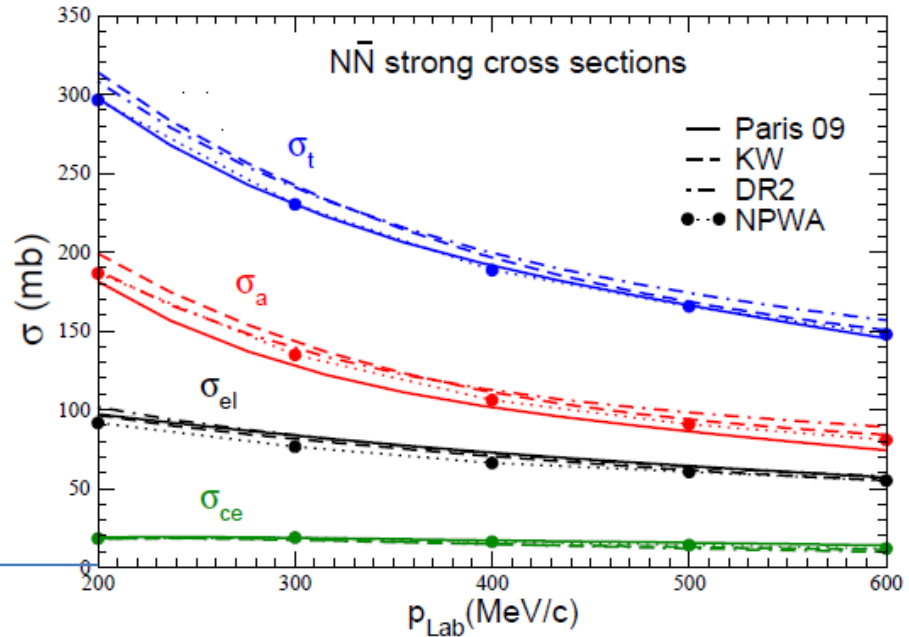


Fig. 1 Integrated strong $\bar{N}N$ cross sections – elastic σ_e (black), annihilation σ_a (red), charge-exchange σ_{ce} (green) and their sum σ_t (blue) – as functions of the \bar{N} laboratory momenta for DR2 (dashed dotted line), KW (dashed line) and Paris 2009 (solid line) optical models. The results of the Nijmegen Partial Wave analysis [7] are indicated by filled circles.

DATA 4000 BUT

TOO MANY PARTIAL WAVES

NO EXCLUSION PRINCIPLE AS IN N-INTERACTION

POSSIBLE BARYONIA

INCONSISTENT MODELS

PARIS : Meson Exchange ,
Dispersion relations

BONN : Chiral expansion
NIMEGHEN Phenomenology

KIOTO-MUNICH Meson exchange

EXAMPLE

S-WAVE BARYONIUM

BES III

P-Pbar

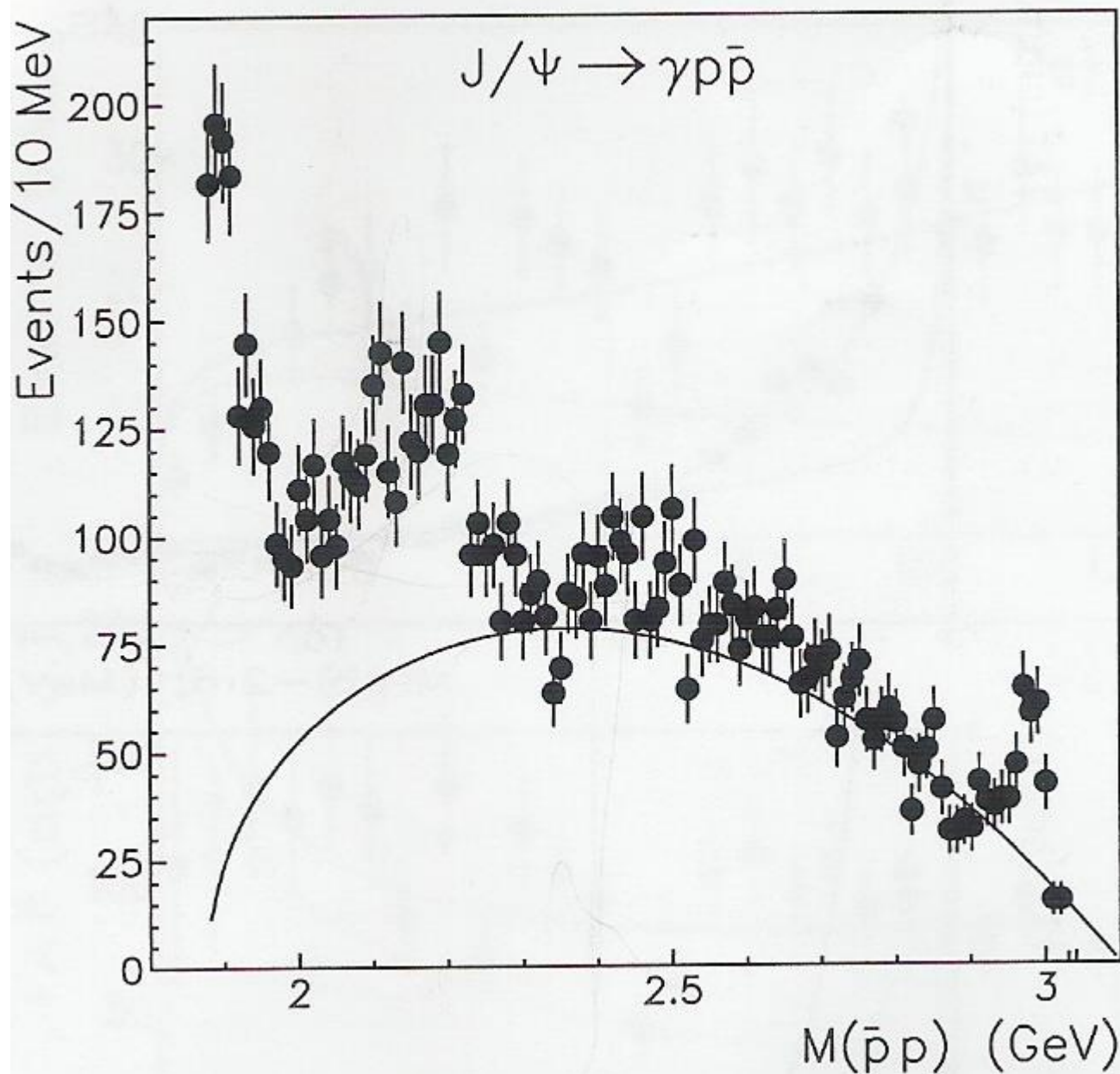
BOUND

STATE

INDICATED

ISOSPIN

UNKNOWN



TESTING S - WAVE AMPLITUDES AT THRESHOLD

Antiprotonic –hydrogen : 1S , 2P levels

- P-Pbar scattering lengths : large differences
 scattering volumes : dramatic differences

state		Exp	Paris 2009	Jülich	KW	DR2
1S_0	$\bar{N}N$		1.02 - i 0.87	0.42 - i 0.91	0.52 - i 0.99	0.65 - i 0.82
	$\bar{p}p$	0.493(92) - i 0.732(146)	0.92 - i 0.67	0.50 - i 0.71	0.57 - i 0.77	0.68 - i 0.64
3S_1	$\bar{N}N$		0.91 - i 0.62	0.93 - i 0.92	1.01 - i 0.79	1.09 - i 0.75
	$\bar{p}p$	0.933(45) - i 0.604(51)	0.82 - i 0.50	0.90 - i 0.74	0.92 - i 0.63	0.98 - i 0.59
S-averaged	$\bar{N}N$		0.94 - i 0.68	0.80 - i 0.92	0.89 - i 0.84	0.98 - i 0.77
	$\bar{p}p$	0.823(57) - i 0.636(75)	0.85 - i 0.54	0.80 - i 0.74	0.83 - i 0.67	0.90 - i 0.60
3P_0	$\bar{N}N$		-3.02 - i 2.50	-0.32 - i 4.01	-3.20 - i 2.28	-2.93 - i 1.83
	$\bar{p}p$	-5.68(123) - i 2.45 (49)	-2.74 - i 2.46	-0.32 - i 3.85	-2.81 - i 1.99	-2.53 - i 1.62

Table 5 Isospin averaged (a_{NN}) and $\bar{p}p$ scattering lengths are compared with those obtained from hydrogen atom level shifts and widths, in fm for S and fm³ for P states. The $\bar{p}p$ values including Coulomb and Δm corrections are taken from [18] for DR2 and KW, from [19] for Paris and from [12] for Jülich model. The statistical averaged value for S-wave is defined as $(^1S_0 + 3\ ^3S_1)/4$ and is given with averaged errors.

CHAOTIC POTENTIALS FOR S-singlet PARTIAL WAVE

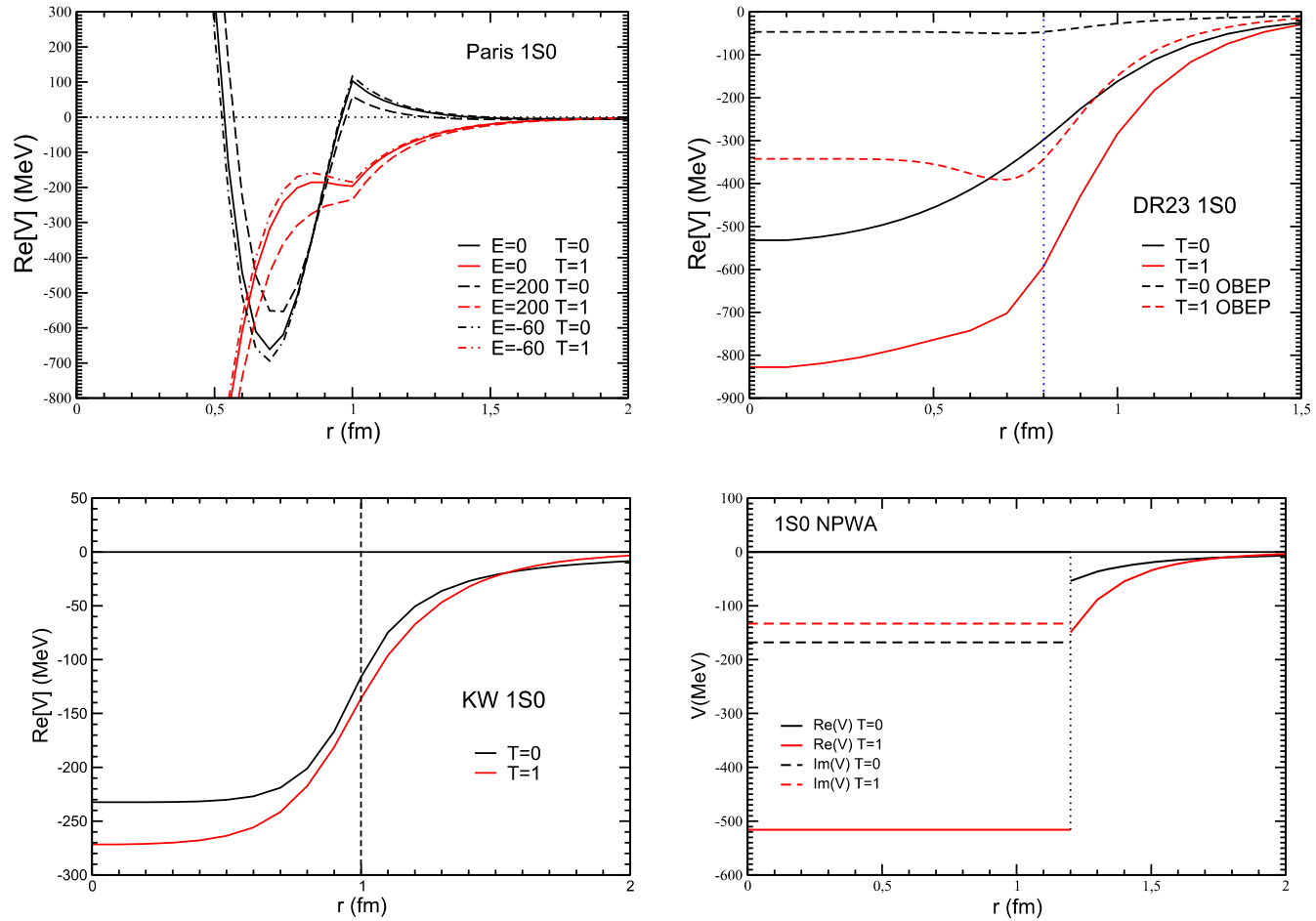


Fig. 15 Real parts of 1S_0 potentials for both isospins (T)

OUR PROGRAM

CHECKING INCONSISTENCIES IN EXISTING N-Nbar THEORIES

J.C.,G.H., SW EUR. Phys. J A

NEW MODEL FOR N-Nbar INTERACTIONS

INCLUDING : SCATTERING DATA = AMPLITUDES IN PHYSICAL REGION

ATOMIC LEVELS = AMPLITUDES FOR NEGATIVE KINETIC ENERGIES

(2) CALCULATE ANTI-PROTONIC NUCLEAR STATES

(3) STUDY SHORT RANGE p-n CORRELATIONS IN NUCLEI WITH PUMA,
a by-product of the experiment

THANK YOU

APPENDICES – if needed

	a_1	r_1	a_1	r_1	a_1	r_1	a_1	r_1
T=0	$^{11}\text{P}_1$		$^{13}\text{P}_0$		$^{13}\text{P}_1$		$^3\text{PF}_2$	
Nijm*	-3.34-1.22i	9.3-1.2i	-3.06-7.23i	-1.7-1.5i	4.36-0.00i	-3.5-0.0i	-	-
Jülich	-2.87-0.36i	-	-2.83-7.82i	-	4.61-0.05i	-	-0.74-1.13i	-
Paris 09	-3.62-0.34i	3.8-0.8i	-8.78-4.99i	0.23-1.1i	5.12-0.02i	-3.4-0.02	-0.49-0.87i	-
KW	-3.36-0.62i	3.7-1.6i	-8.83-4.45i	0.25-0.97i	4.73-0.08i	-3.5-0.1i	-0.46-1.09i	-
DR2	-3.28-0.78i	4.2-2.3i	-8.53-3.50i	0.63-1.0i	5.14-0.09i	-3.4-0.1i	-0.59-0.85i	-
T=1	$^{31}\text{P}_1$		$^{33}\text{P}_0$		$^{33}\text{P}_1$		$^3\text{PF}_2$	
Nijm*	0.66-0.18i	3.3-20i	2.33-0.92i	-10-0.7i	-2.02-0.70i	4.7-2.8i	-	-
Jülich	0.80-0.34i	-	2.18-0.19i	-	-2.04-0.55i	-	-0.48-0.34i	-
Paris 09	1.00-0.77i	-3.7-9.8i	2.74-0.00i	-5.2-0.01i	0.28-4.11i	-3.0-2.0i	-0.13-0.21i	-
KW	0.71-0.47i	-8.3-21i	2.43-0.11i	-5.8-0.43i	-2.17-0.95i	2.7-3.5i	-0.30-0.45i	-
DR2	1.02-0.43i	-11-10i	2.67-0.15i	-5.4-0.53i	-2.02-0.70i	4.6-3.9i	-0.04-0.53i	-

Table 3 P waves $\bar{N}N$ low energy parameters (in fm^3) for the considered optical models: Jülich results are taken from Tab 3 of Ref. [12], KW and DR2 from [18], Paris 2009 have been recomputed and are in agreement with [44]. The values of Nijmegen are obtained by extrapolating the phase shifts from Figures 2 and 3.

CHOICE OF PARAMETERS TO DESCRIBE FINAL MESON INTERACTIONS and P(Q)

$$p \bar{p} \rightarrow Q_{ini}=0 ; \quad n \bar{p} \rightarrow Q_{ini}=-1 \quad \text{PARAMETER}$$

$$\begin{array}{lll} \pi (+) \text{ NN} & \rightarrow & \text{NN} & Q \rightarrow Q-1 \\ \pi (+) \text{ n} & \rightarrow & \pi (0) \text{ p} & Q \rightarrow Q-1 & \omega(+)$$

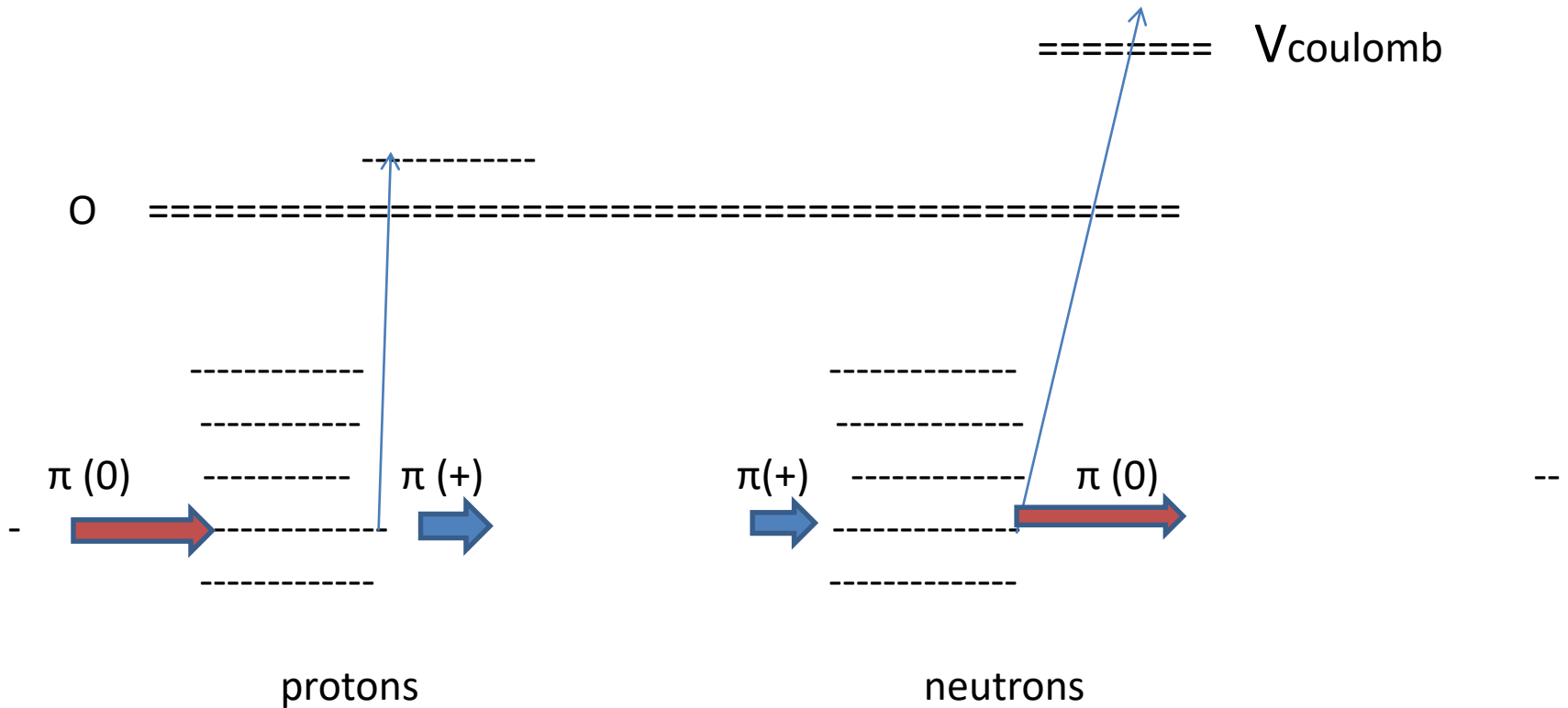
$$\begin{array}{lll} \pi (-) \text{ NN} & \rightarrow & \text{NN} & Q \rightarrow Q+1 \\ \pi (-) \text{ p} & \rightarrow & \pi (0) \text{ n} & Q \rightarrow Q+1 & \omega(-)$$

$$\begin{array}{lll} \pi (0) \text{ n} & \rightarrow & \pi (-) \text{ p} \\ \pi (0) \text{ p} & \rightarrow & \pi (+) \text{ n} \end{array} \quad \begin{array}{l} \text{different} \\ \downarrow \end{array} \quad \begin{array}{ll} Q \rightarrow Q-1 & \lambda(-) \\ Q \rightarrow Q+1 & \lambda(+)$$

$$\pi (0) \rightarrow \text{lost} \quad \omega(0)$$

$$\omega \sim 0.1-0.2 \quad ; \quad \lambda \sim 0.15 - 0.40 \quad \text{from data}$$

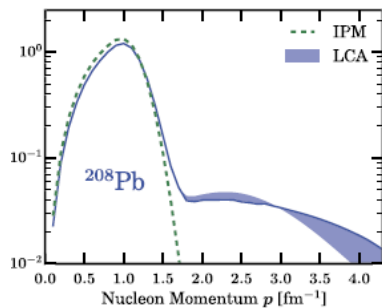
PAULI BLOCKING depends on nucleon binding **and momentum**



Charge exchange differs from its inverse due to exclusion and Coulomb barrier

Difference depends on nucleon momenta.

Nucleon momenta in a nucleus
 Fermi gas sector and p-n short range correlations sector.
 M Duer+ Phys Rev Let 112 J. Lab electron scattering



Ryckebusch +
 Phys L. B792

J. Ryckebusch et al. / Physics I

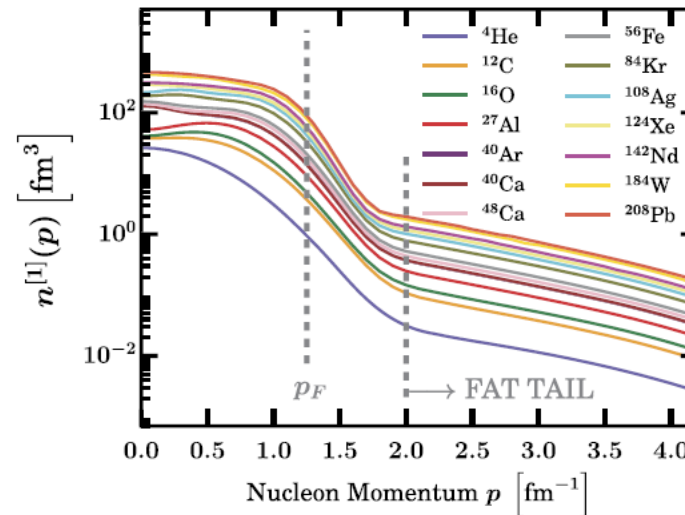


Fig. 2. The momentum distribution for 14 nuclei across the nuclear mass table. The $n^{[1]}(p)$ are computed in LCA with a "hard" central correlation function g_c adopting the normalization convention $\int dp p^2 n^{[1]}(p) = A$.

Initiated by Campi and Bouysy, old
 problem of correlations revived with
 different physics

SHORT LIFE of ATOMS

Radii = $57 / Z n^2$ fm

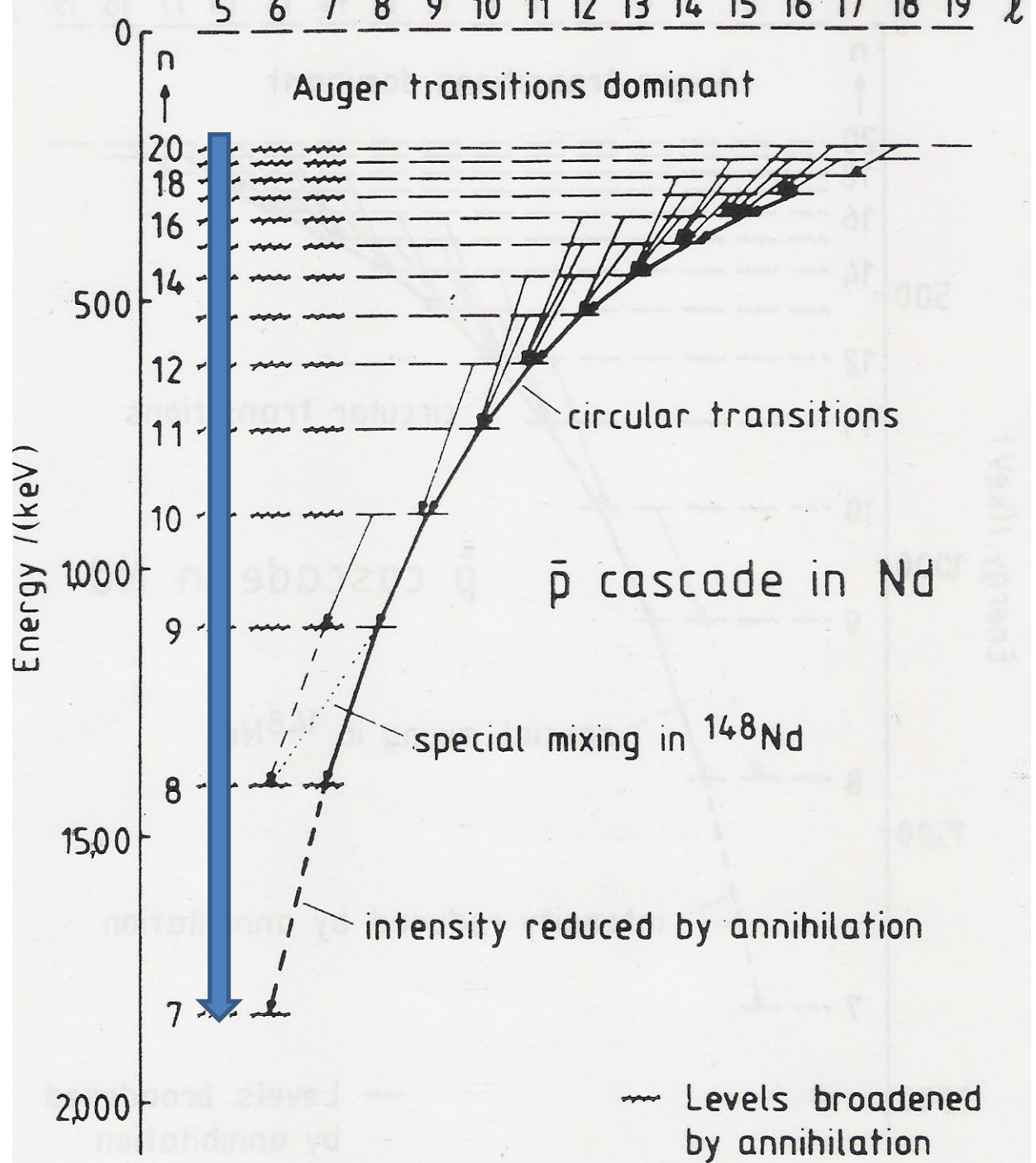
High l levels
 $\Psi / r^l \sim \text{const}$
inside nuclei

FINAL
PRODUCTS

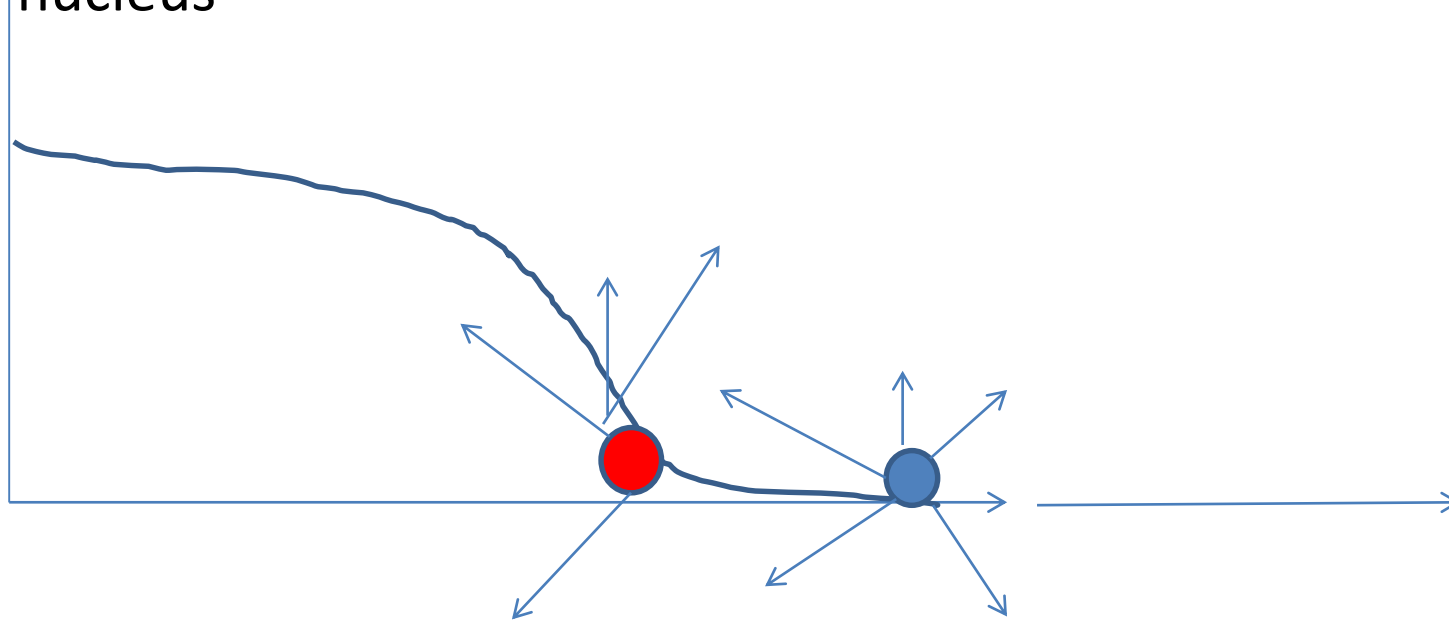
X-rays

Nuclei

Pions



Antiproton ● a 2000 MeV bomb on the side of nucleus



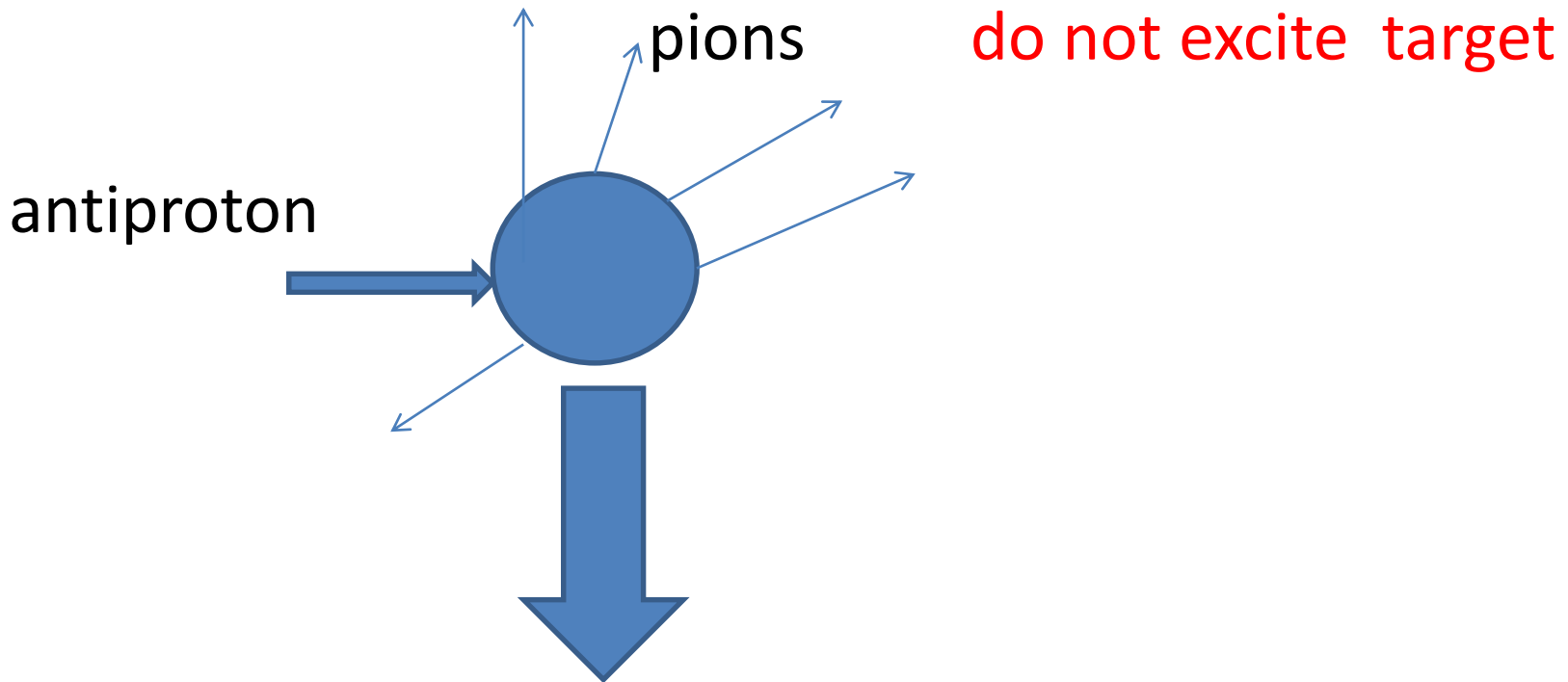
→ Mesons from
NN-bar → $\pi \pi \pi$

● Nucleus destroyed, pions detected, less peripheral

● Cold residual nuclei detected, more peripheral

X-rays regions in between

Radiochemical measurements of final non excited nuclei Munich – Warsaw /CERN



A-1 NUCLEI MEASURED \rightarrow RATIO $(N-1)/(Z-1)$

DETERMINATION OF CAPTURE ORBIT via $(A-1)/\text{TOTAL}$

ANALYSIS OF COLD CAPTURES

$$\frac{\sigma(N-1)}{\sigma(Z-1)} = \frac{N P_{\text{emission}N}}{Z P_{\text{emission}Z}} R_{n/p} f_{\text{HALO}}$$

$R_{n/p}$ relative rate of absorptions (p-bar n) / (p-bar p)

P_{emission} chance for mesons not to excite the nucleus ~10%

Result f_{HALO} excess of neutrons in the capture region
 estimated from $\sigma(A-1) / \sigma(\text{total})$

Presentation : if capture region is known
 => $R_n - R_p =$ difference of Rms radii is calculated

Excess of neutrons
over protons
Reduced by N/Z

Lubinski PRC 57
Munich Warsaw

With known
capture orbit

$R_{\text{ms}}(n) - R_{\text{ms}}(p)$
Extracted

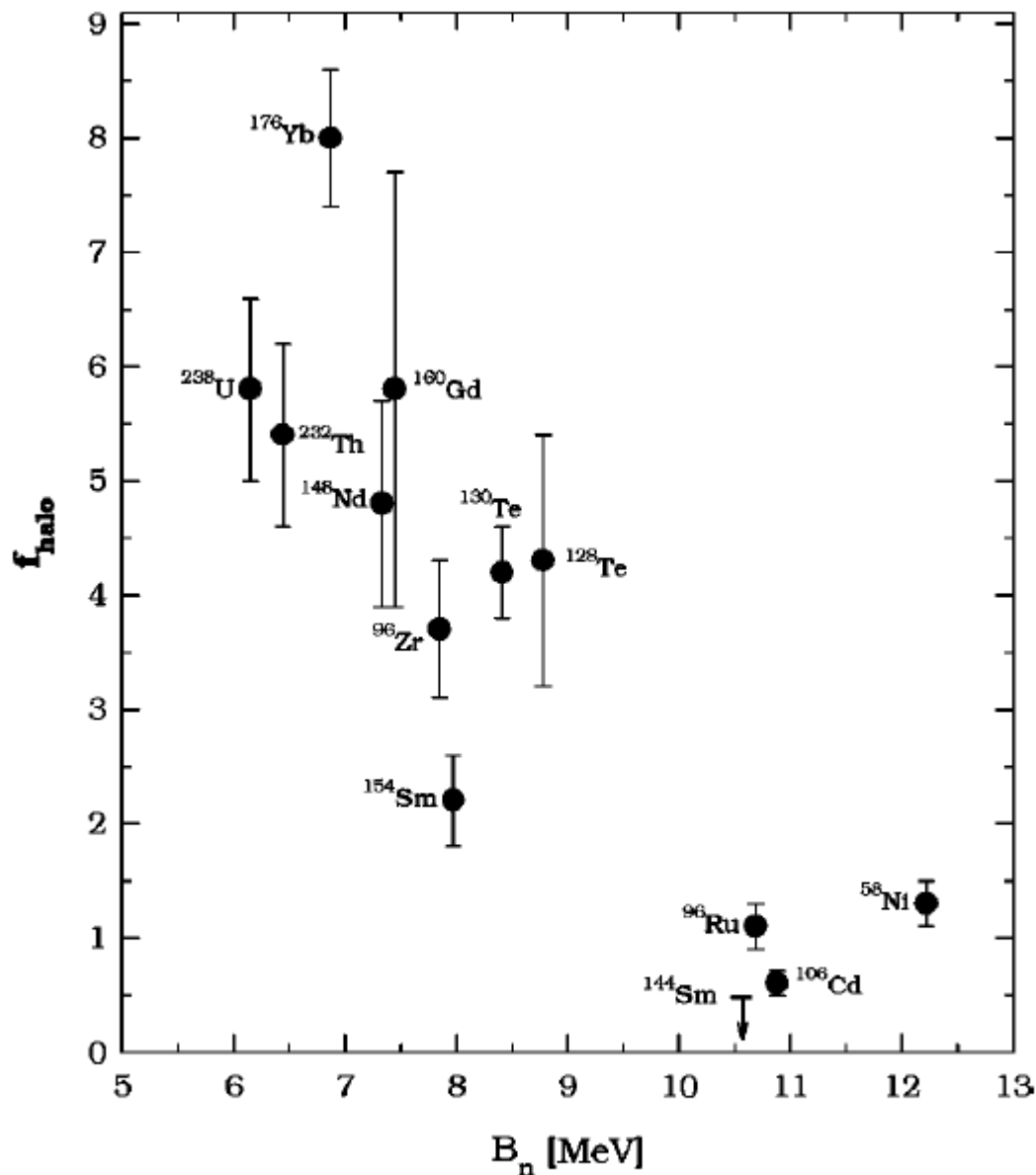


FIG. 3. Neutron halo factor (defined in the text) as a function of the target neutron separation energy B_n .

Nitrogen , Riedlberger + PRev C40 (1989) High statistics , No hydrogen contamination, magnetic spectrometer

: Experimental,[21], and fitted charge multiplicities $P[Q]$ in Nitrogen .

Q	exp	fit
3	1.2(.2)	0.28
+2	3.9(.4)	2.25
+1	14.2(.8)	15.6
0	39.5(1.0)	40.1
-1	31.1(.8)	32.1
-2	8.0(.5)	8.5
-3	2.1(.3)	0.44
$\langle n^\pm \rangle$	2.89(8)	2.91(0.05)
χ^2		7.5

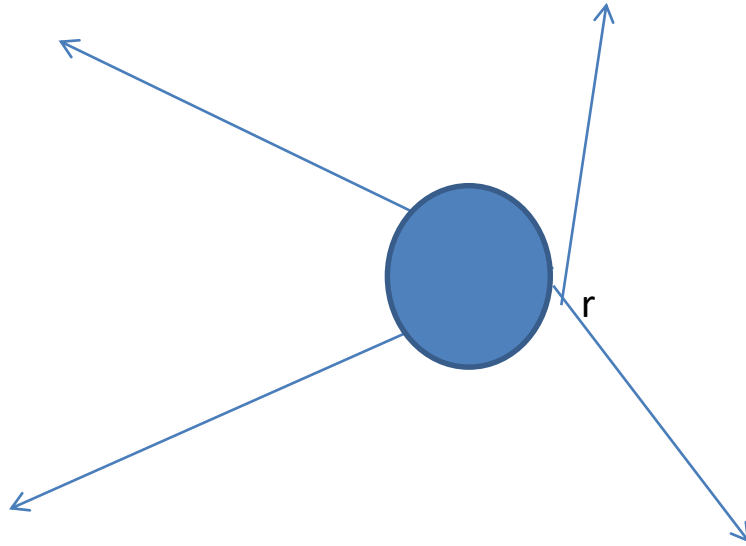
$$R_{n/p} \cdot f^h = 0.77(.04).$$

$$\omega^+ = 0.16 ; \omega^- = .17 ; \lambda^+ = .16 ; \lambda^- = 0.10$$

END POINTS INDICATE DOUBLE PION CHARGE
EXCHANGE ON RESIDUAL Carbon = $\alpha\alpha\alpha$

CALCULATION OF PARAMETERS

MESONS ARE FAST average momenta ~ 400 MeV/c \rightarrow eikonal approximation



$$T_{\text{expt}}(\mathbf{r}, \mathbf{k}) = \exp \left[-\lambda_{\text{expt}} \int_0^{\infty} ds \rho_p(\mathbf{r} - s\hat{\mathbf{k}}) \right].$$

Survival amplitude $T = 1 - \omega$

Average over momenta , directions
number of mesons

Charge exchange $\lambda = \sigma * \text{Pauli Blocking factor}$
NN absorption $\rho \rightarrow \rho\rho$

mostly surface
mostly centre

FERMI MOMENTUM AT NUCLEAR SURFACE ?

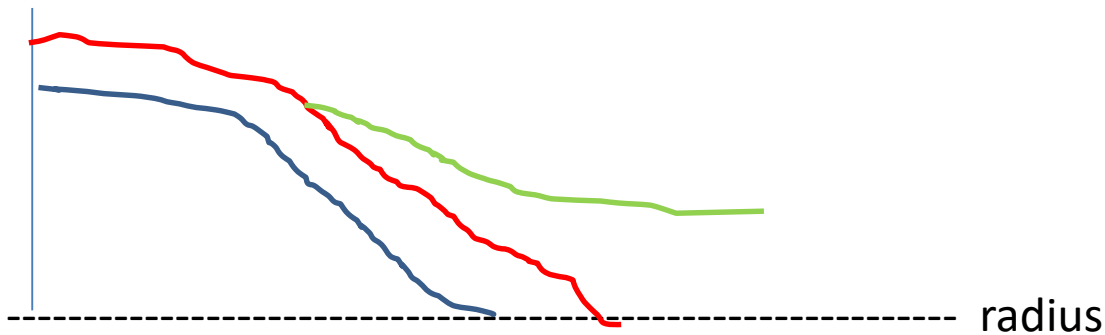
Fermi gas $K_{\text{FERMI}} \sim \rho^{1/3}$

$$\rho(x, x') = \sum \phi(x) \phi(x')^*$$

Wigner function

$$= \rho(x/2 + x'/2) j_1(K_{\text{FERMI}} |x - x'|)$$

correlation function



$K_{\text{FERMI}}(r)$ Fermi gas :

K_{fermi} shellmodel

X Campi, A Bouyssy , 1973

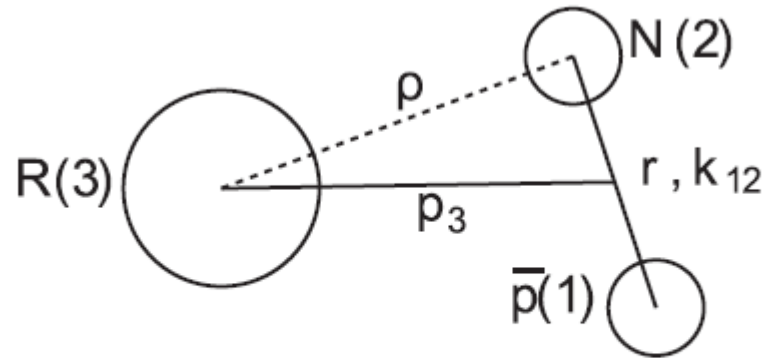


FIG. 1. Quasi-three-body system: (1) antiproton, (2) nucleon, and (3) residual system. Jacobi coordinates: momentum p_3, k_{12} and space ρ, r .

In atoms

Kinetic N-Nbar ENERGY in CM system is **negative**

$$E_{\text{CM}} = 2M - \text{Binding} - \text{Recoil}$$

$\bar{N} - N$ quasi-bound states