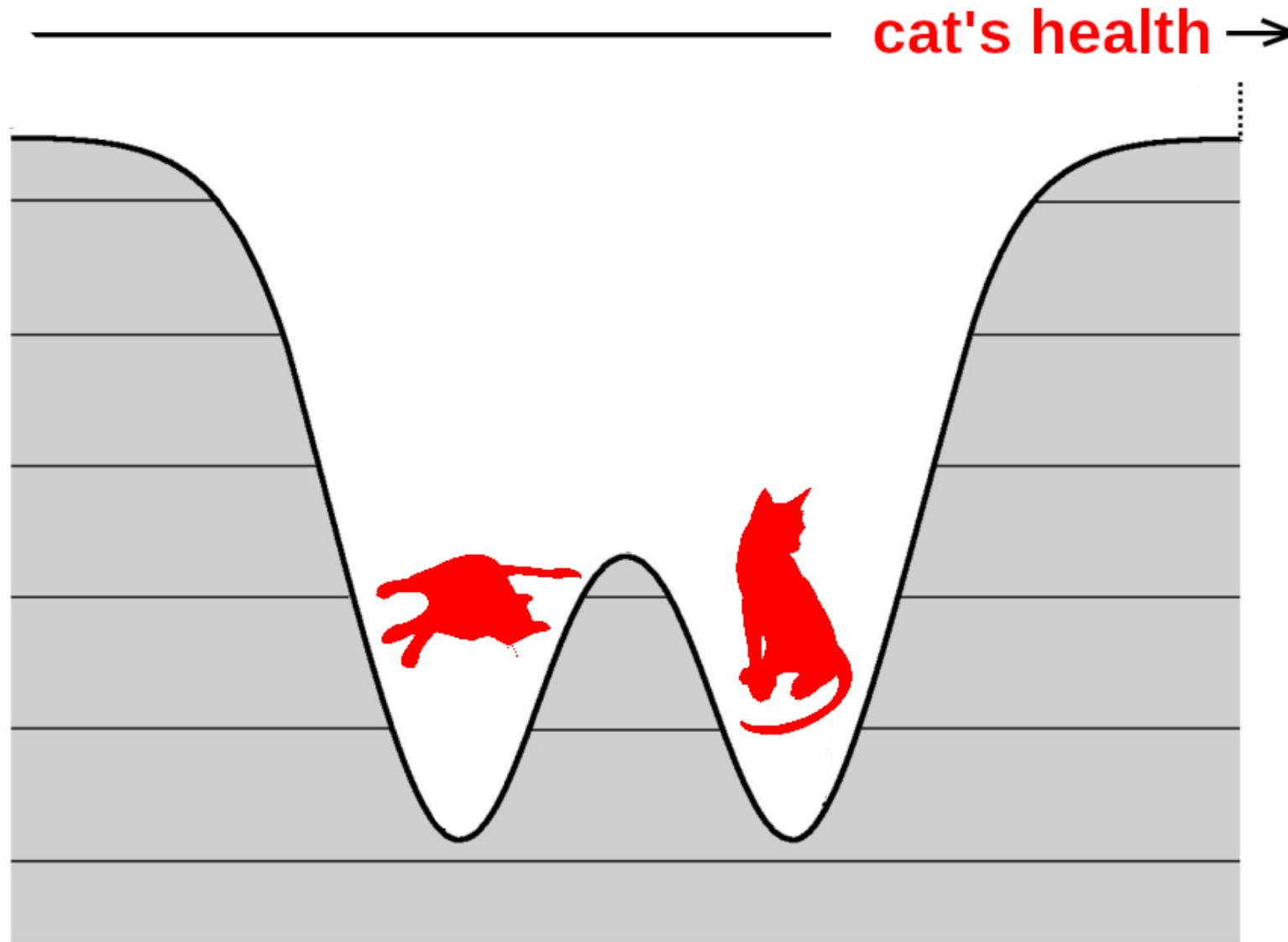


# First direct measurement of nuclear chirality in $^{128}\text{Cs}$ .

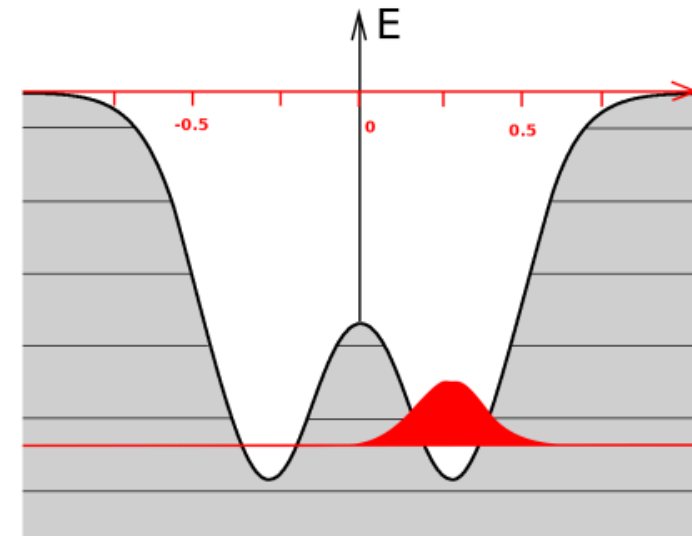
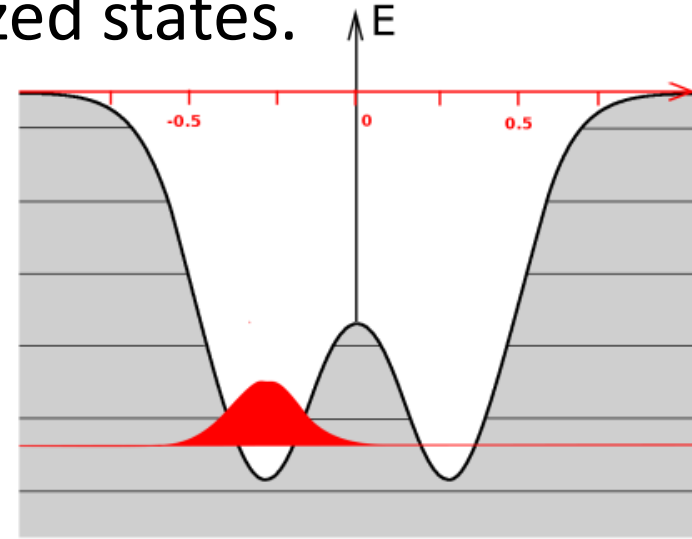
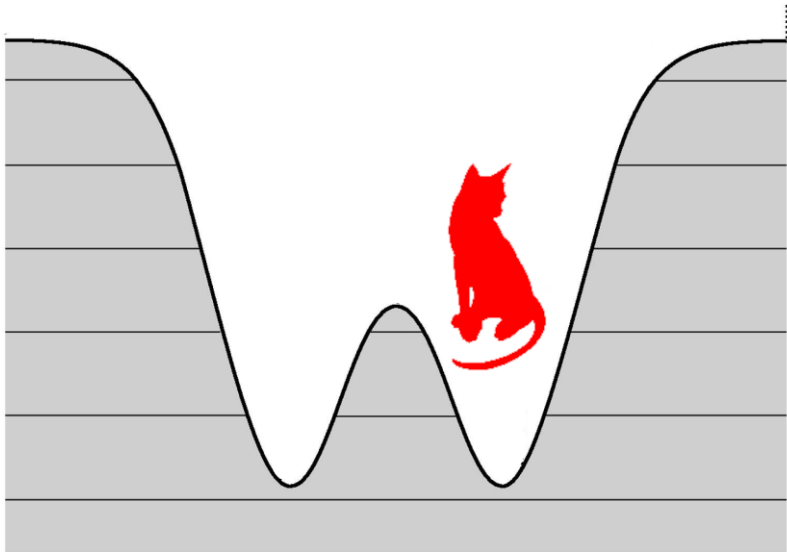
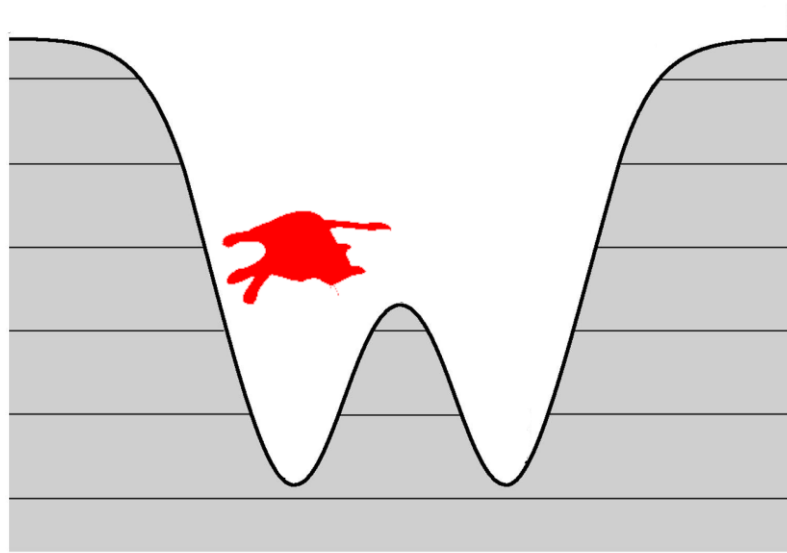
Ernest Grodner

BP1

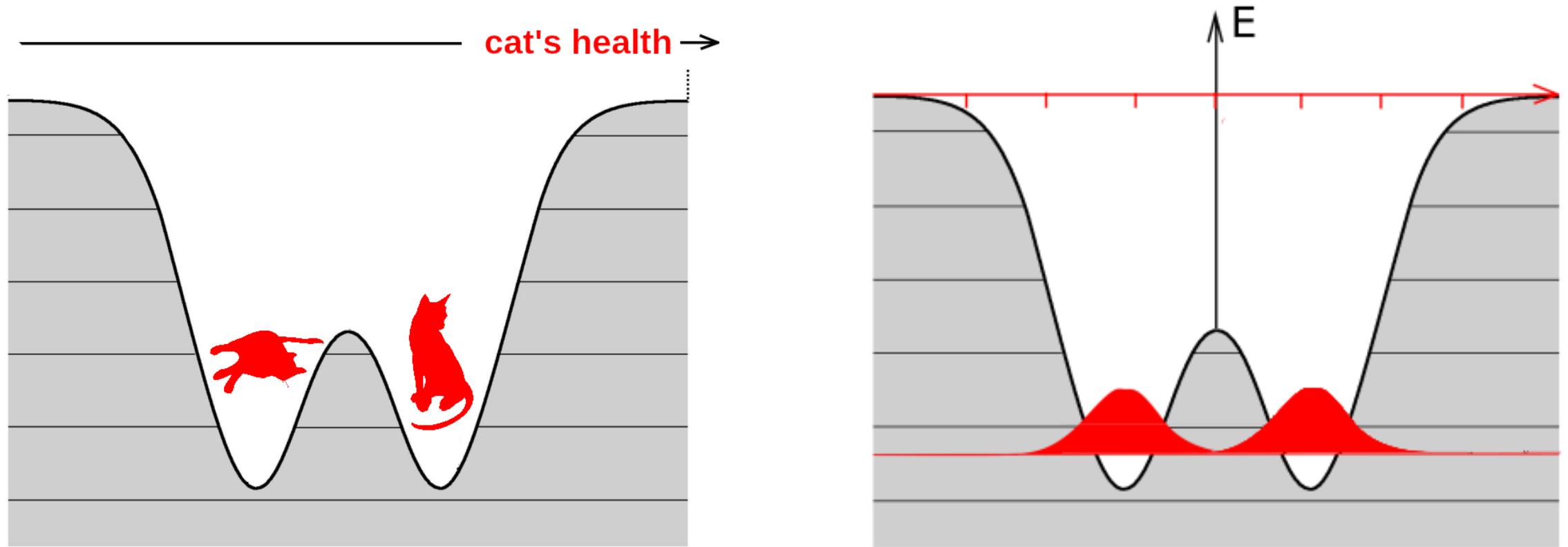
Nuclear chirality is a today's nuclear spectroscopy  
Schrodinger's cat box problem.



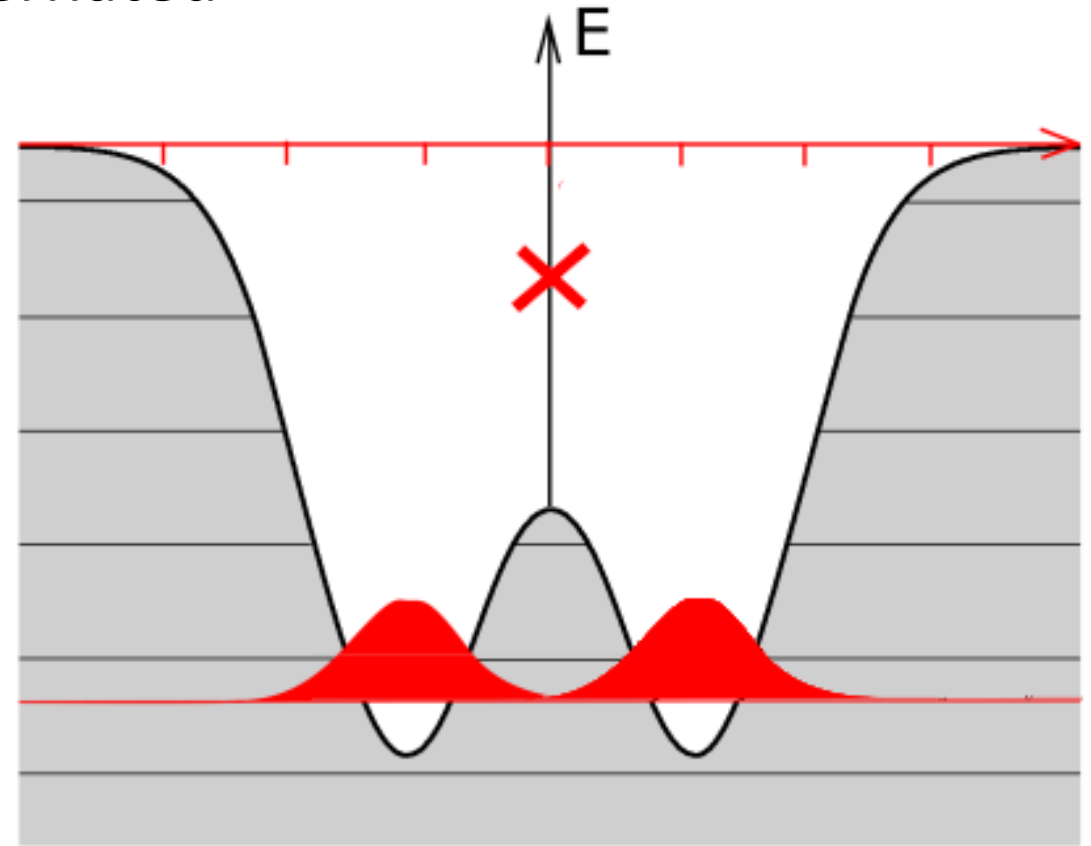
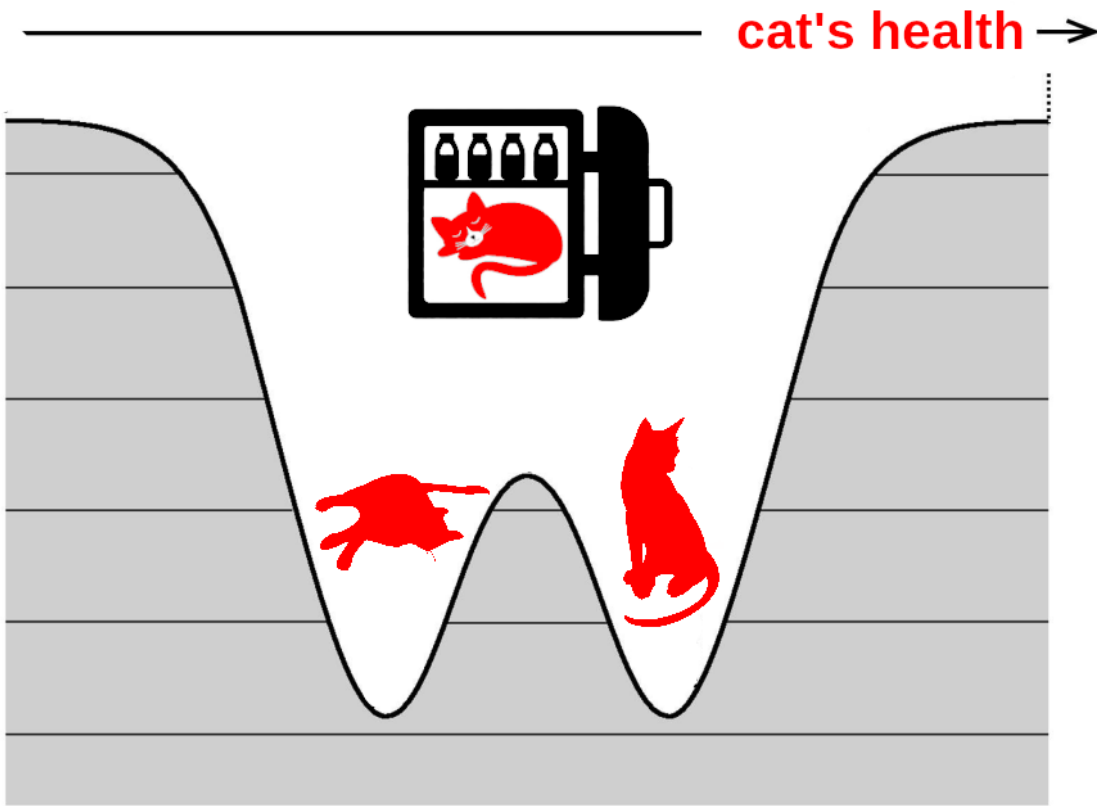
Taking a look into the cat's box allowed, then cats state can be measured. We deal with localized states.



Taking a look into the cat's box forbidden.  
We must deal with superimposed states



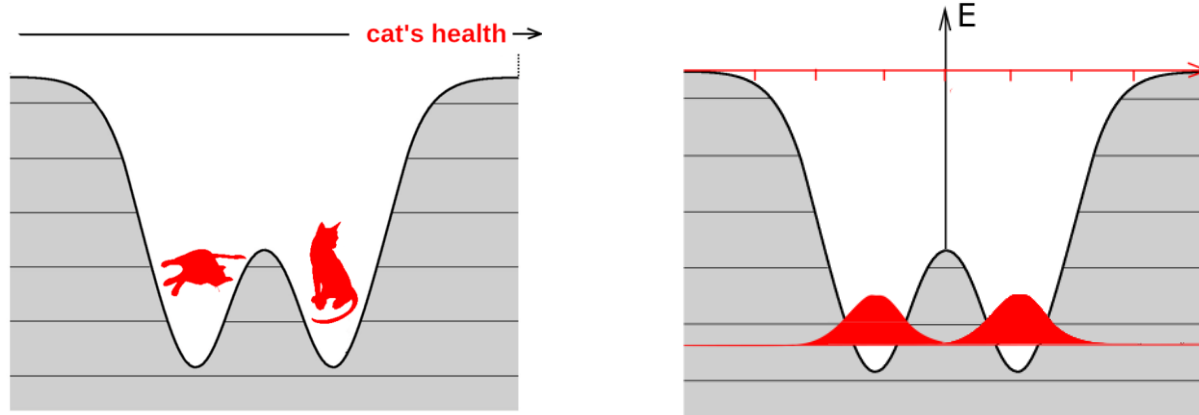
Cat's health expectation value on superimposed states?  
Neither alive nor dead, rather = hibernated



Attention!

Now the clue of  
the lecture

Superimposed states of a cat in the box  
**Symmetry breaking cat inside**

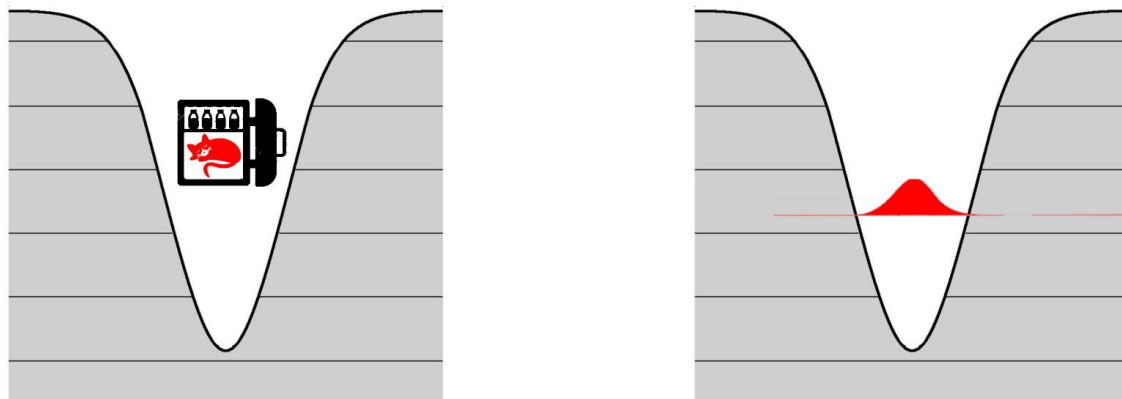


Measured cat's health:  
hibernated

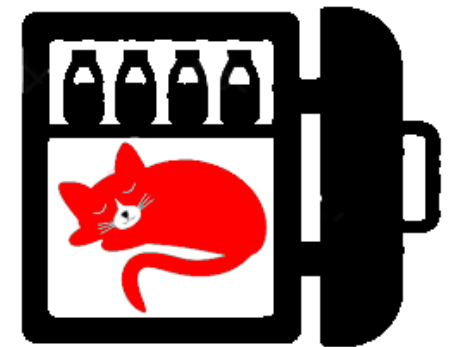


But what if we put a hibernated cat in the box  
in a first place?

**Symmetry conserved cat inside**

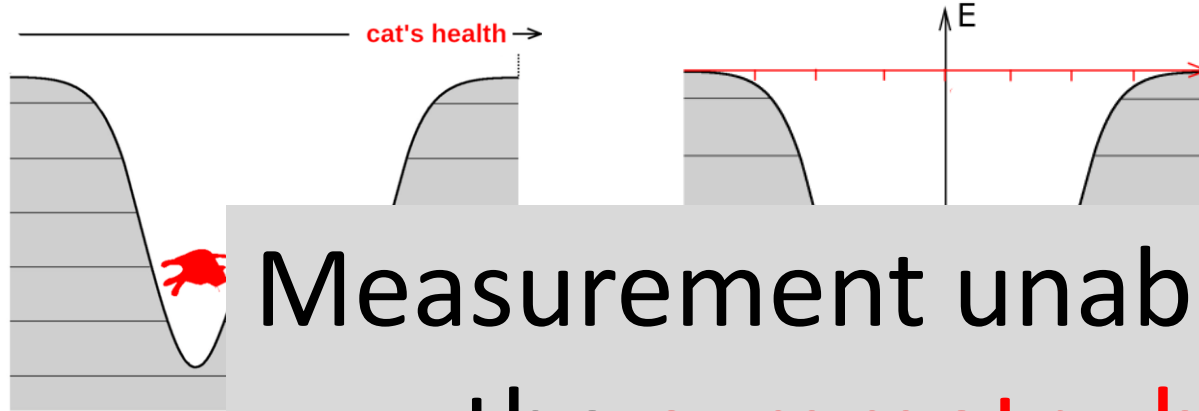


Measured cat's health:  
hibernated



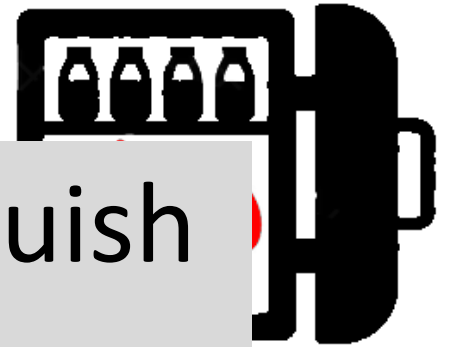
Superimposed states of a cat in the box  
**Symmetry breaking cat inside**

Measured cat's health:  
hibernated

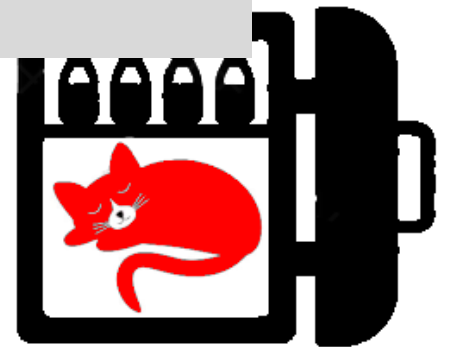
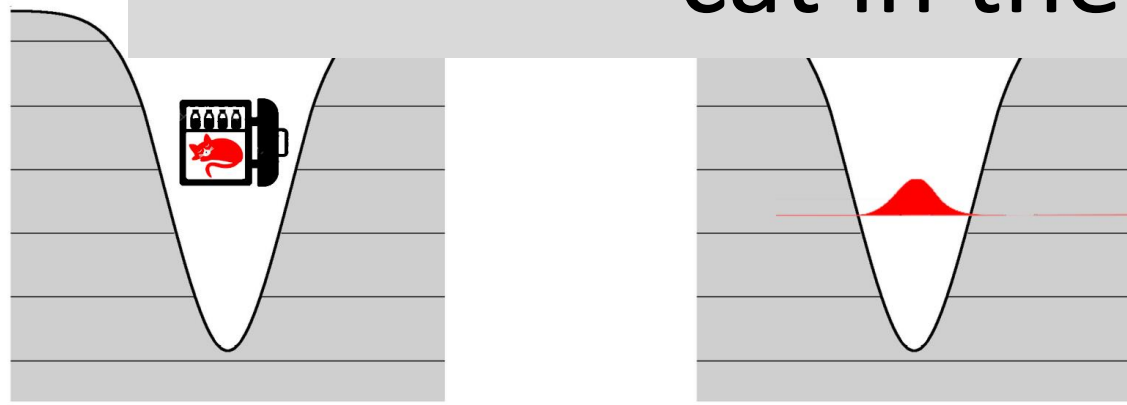


But wh  
in a fir  
**Symme**

Measurement unable to distinguish  
the **symmetry broken** from  
The **symmetry conserved**  
cat in the box.



's health:  
ted





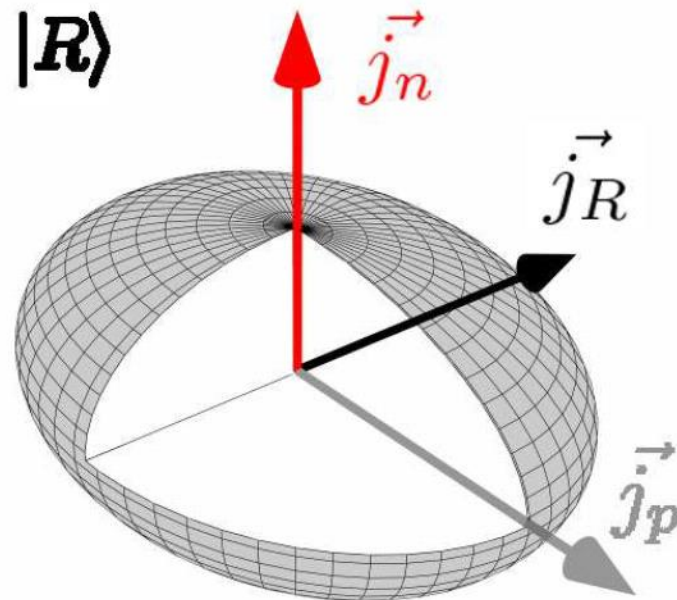
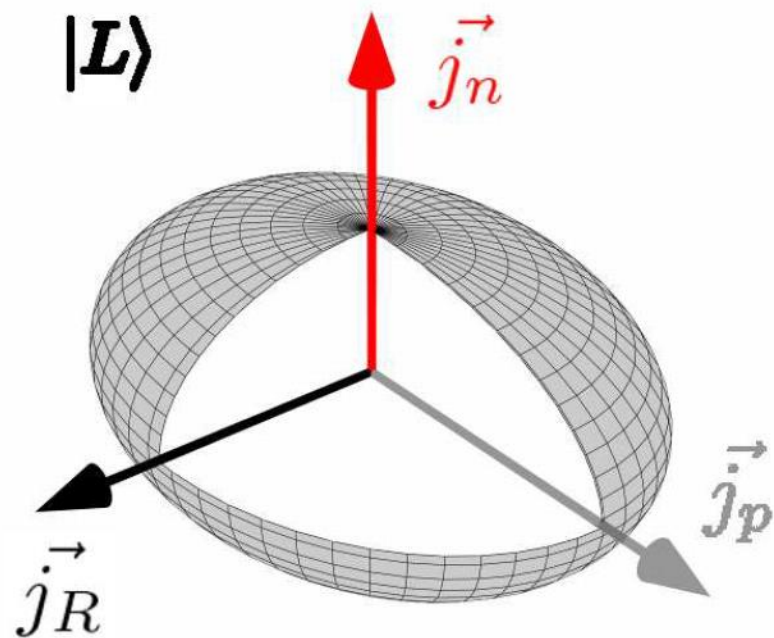
## Nuclear chirality

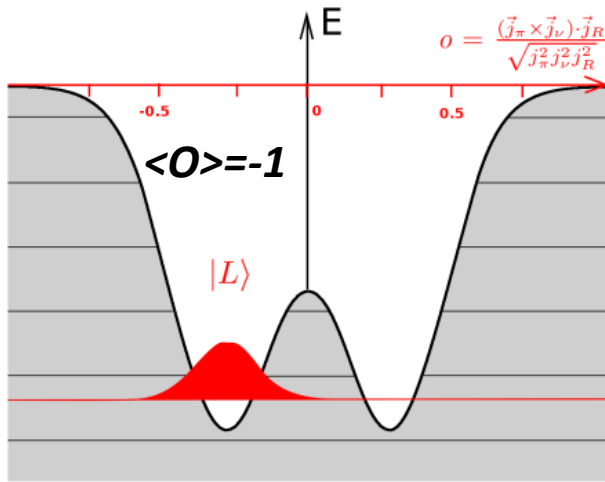
Handedness instead  
of cat's health parameter

$$O = \frac{(\vec{j}_\pi \times \vec{j}_\nu) \cdot \vec{j}_R}{\sqrt{j_\pi^2 j_\nu^2 j_R^2}}$$

$$\langle O \rangle = -1$$

$$\langle O \rangle = +1$$



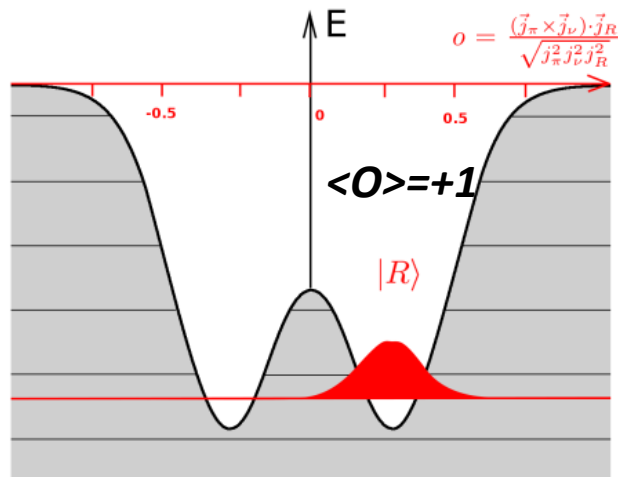


**Fusion-evaporation reactions used to produce highly excited odd-odd isotopes at HIL (U200p cyclotron).**

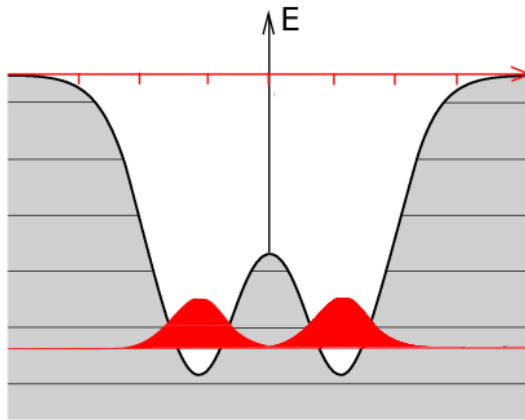
**A nucleus cools-down emitting particles and gamma quanta.**

**At some point it must choose spontaneously the handedness.**

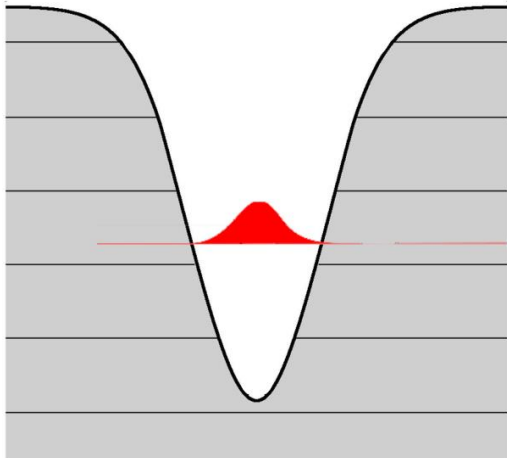
**Spontaneous chiral symmetry breaking in nuclear physics.**



Broken symmetry



Conserved symmetry



In both cases the same measured handedness value = 0.0  
(a hibernated cat again)

$$\langle O \rangle = \frac{(\vec{j}_\pi \times \vec{j}_\nu) \cdot \vec{j}_R}{\sqrt{j_\pi^2 j_\nu^2 j_R^2}} = 0.0$$

Expectation value of handedness  
does not distinguish the symmetry broken  
from the symmetry conserved nucleus  
since handedness is a signed value.

Measuring unsigned observable may distinguish symmetry broken from symmetry conserved state.

$$\langle g \rangle = +1$$

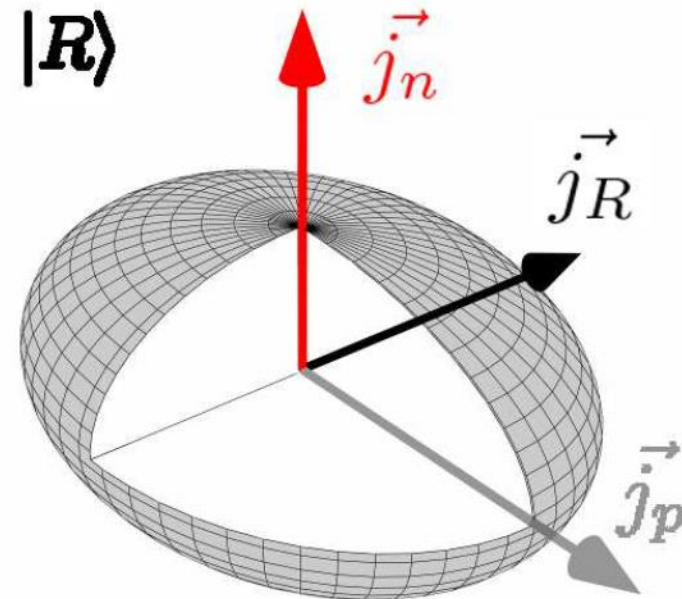
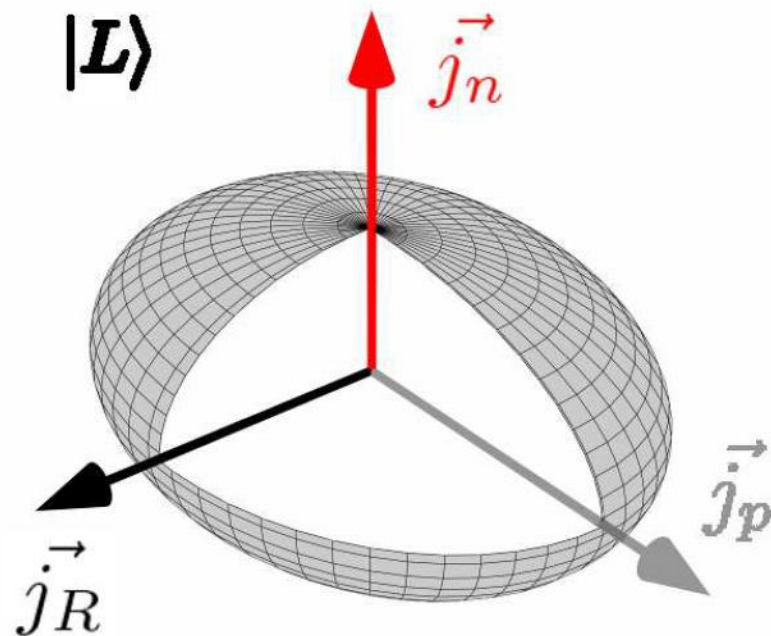
Left handed

$$\langle g \rangle = 0$$

planar

$$\langle g \rangle = +1$$

Right handed



Magnetic dipole moment is a hit! Measured value: the g-factor

$$\frac{1}{\langle J^2 \rangle} \left( g_p \langle \vec{j}_n \cdot \vec{j}_R \rangle + g_n \langle \vec{j}_p \cdot \vec{j}_R \rangle + g_R \langle \vec{j}_p \cdot \vec{j}_n \rangle \right)$$

$$\langle g \rangle = 0$$

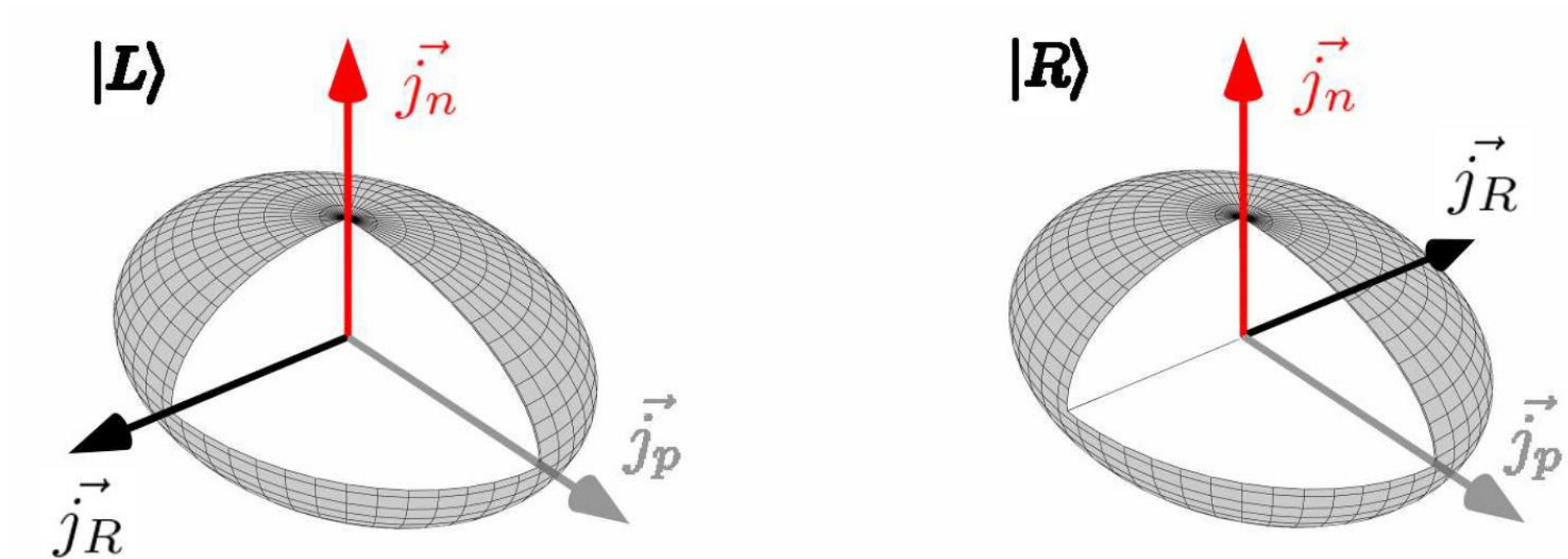
Left handed

$$\langle g \rangle = +0.1$$

planar

$$\langle g \rangle = 0$$

Right handed



**First Measurement of the  $g$  Factor in the Chiral Band:  
The Case of the  $^{128}\text{Cs}$  Isomeric State**

E. Grodner,<sup>1,2</sup> J. Srebrny,<sup>3</sup> Ch. Droste,<sup>2</sup> L. Próchniak,<sup>3</sup> S. G. Rohoziński,<sup>2</sup> M. Kowalczyk,<sup>3</sup> M. Ionescu-Bujor,<sup>4</sup> C. A. Ur,<sup>5</sup>  
K. Starosta,<sup>6</sup> T. Ahn,<sup>7</sup> M. Kisieliński,<sup>3</sup> T. Marchlewski,<sup>3</sup> S. Aydin,<sup>8,10</sup> F. Recchia,<sup>9</sup> G. Georgiev,<sup>11</sup> R. Lozeva,<sup>11</sup> E. Fiori,<sup>11</sup>  
M. Zielińska,<sup>3</sup> Q. B. Chen,<sup>12</sup> S. Q. Zhang,<sup>12</sup> L. F. Yu,<sup>12</sup> P. W. Zhao,<sup>12</sup> and J. Meng<sup>12,13</sup>

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## Examination of nuclear chirality with a magnetic moment measurement of the $I = 9$ isomeric state in $^{128}\text{Cs}$

E. Grodner, M. Kowalczyk, M. Kisieliński, J. Srebrny, L. Próchniak, Ch. Droste, S. G. Rohoziński, Q. B. Chen, M. Ionescu-Bujor, C. A. Ur, F. Recchia, J. Meng, S. Q. Zhang, P. W. Zhao, G. Georgiev, R. Lozeva, E. Fiori, S. Aydin, and A. Nalęcz-Jawecki

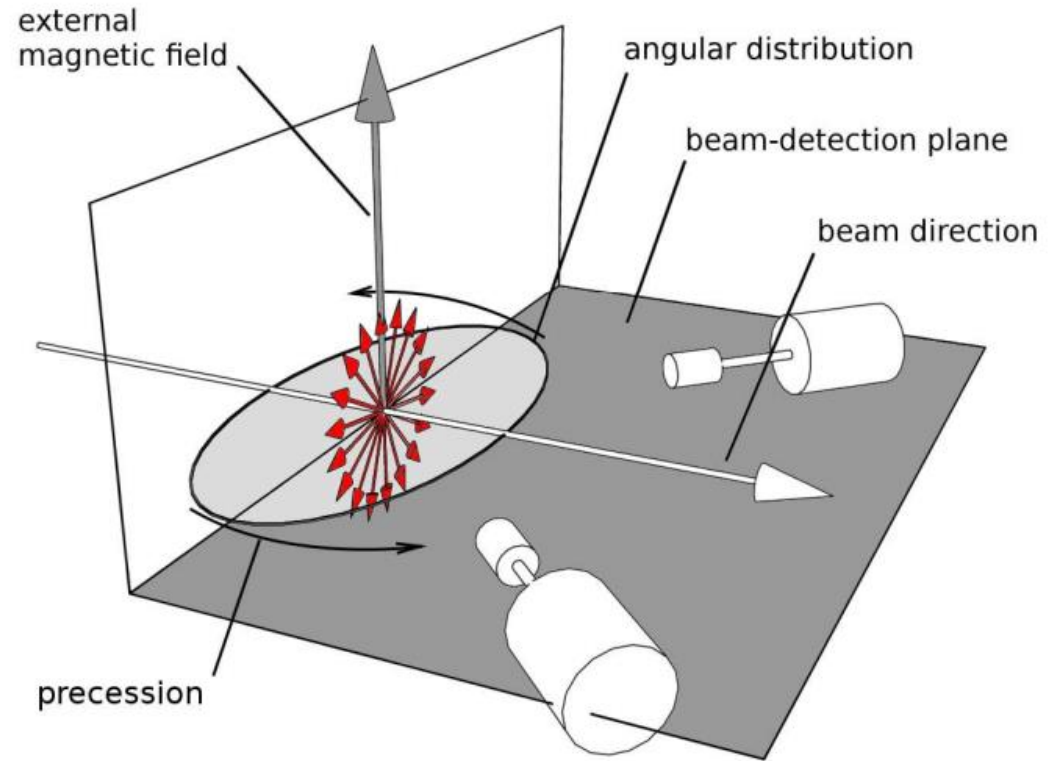
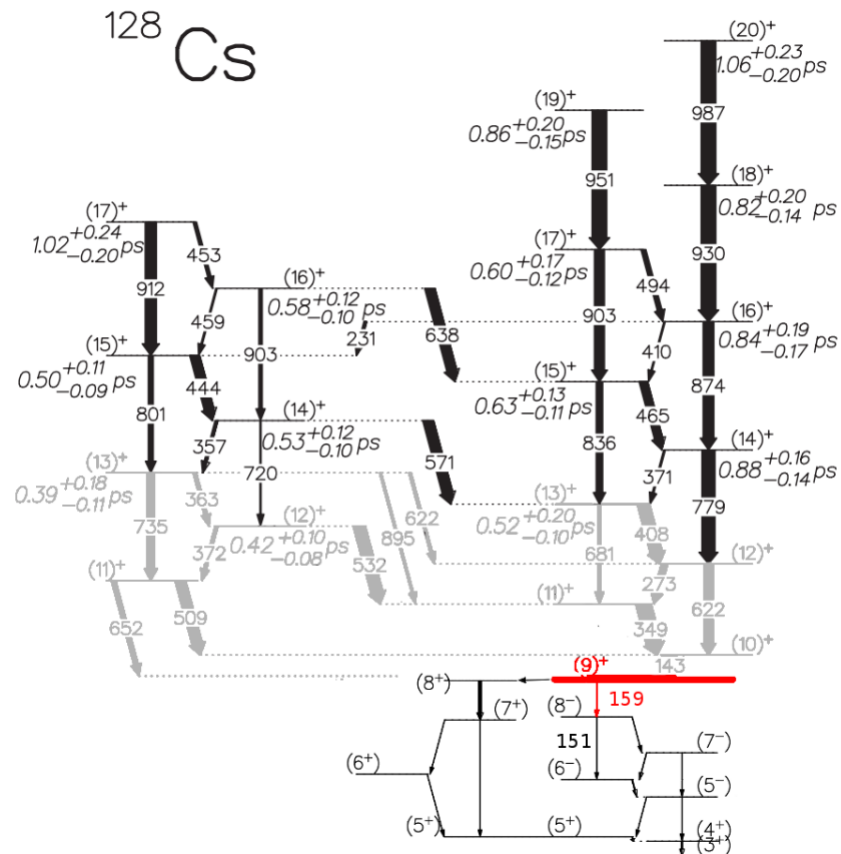
Phys. Rev. C **106**, 014318 – Published 28 July 2022

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### ABSTRACT

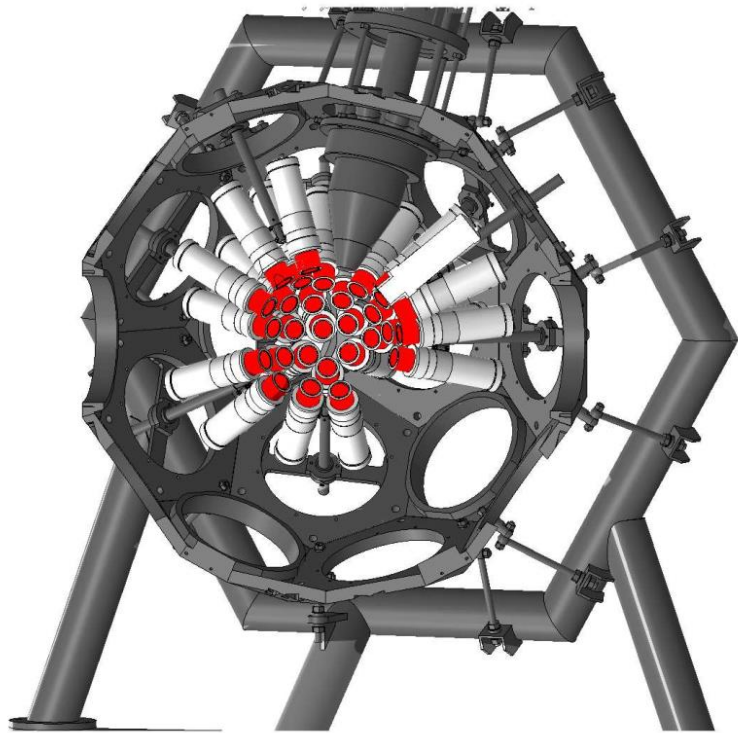
The  $g$  factor of the isomeric  $I = 9^+$  bandhead of the yrast states in  $^{128}\text{Cs}$  is obtained from the time differential perturbed angular distribution measurement performed with the electromagnet at IPN Orsay. An external magnetic field of 2.146 T at the target position was attained with GAMPE reaction chamber surrounded by four high-purity germanium detectors, of which two were low-energy photon spectrometer type. The results are in accordance with  $\pi h_{11/2} \otimes \nu h_{11/2}^{-1}$   $I = 9^+$  bandhead assignment and are discussed in the context of chiral interpretation of the  $^{128}\text{Cs}$  nucleus as a composition of the odd proton, odd neutron, and even-even core with their angular momentum vectors. The obtained  $g$ -factor value was compared with predictions of the particle-rotor model. The experimental  $g$  factor corresponds to the nonchiral geometry of the isomeric bandhead. This observation indicates the existence of the chiral critical frequency in  $^{128}\text{Cs}$  and may explain the absence of the chiral doublet members for  $I < 13\hbar$ .

Just two detectors with magnet on a table. The cheapest experiment with an expensive idea.

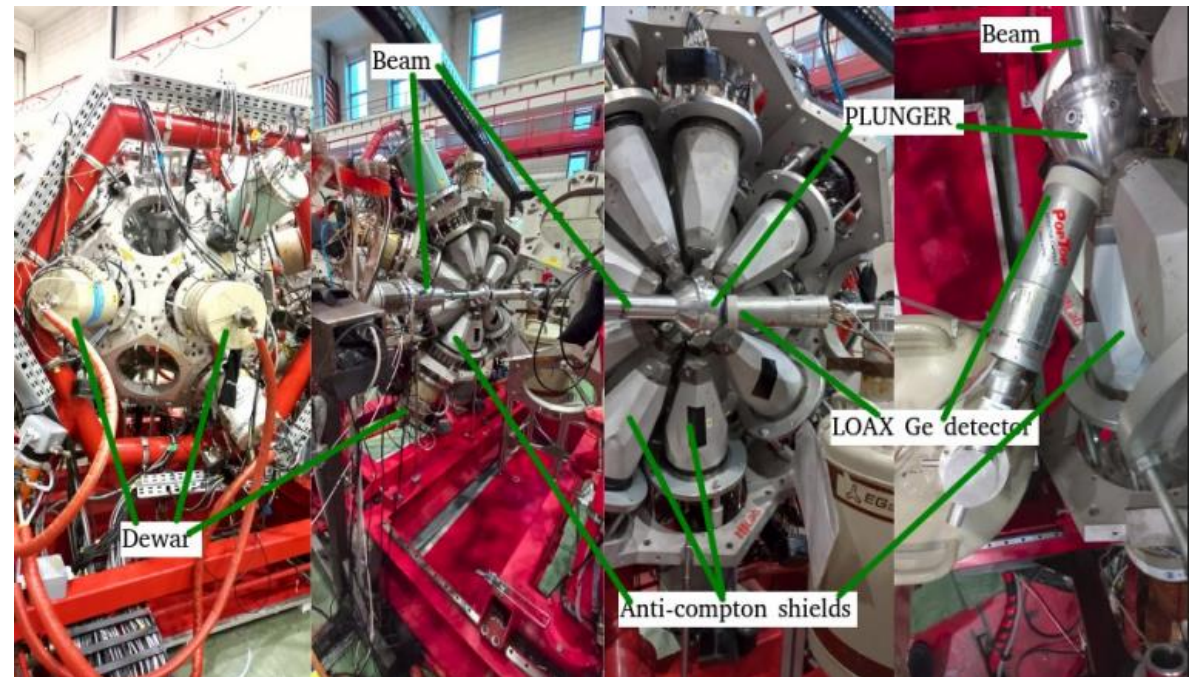




Future: preparations for similar measurements in other excited states. Fast-Timing and Plunger lifetime measurements.  
PHD thesis of Adam Nałęcz-Jawecki (NCNR).



EAGLE-EYE



EAGLE-PLUNGER