

8th row: 119 (ongoing exp. at RIKEN), 120 (ongoing exp. at DUBNA)

Superheavy elements (Transactinides)

SHE = TAN $(Z \ge 104)$

15 elements discovered in 50 years

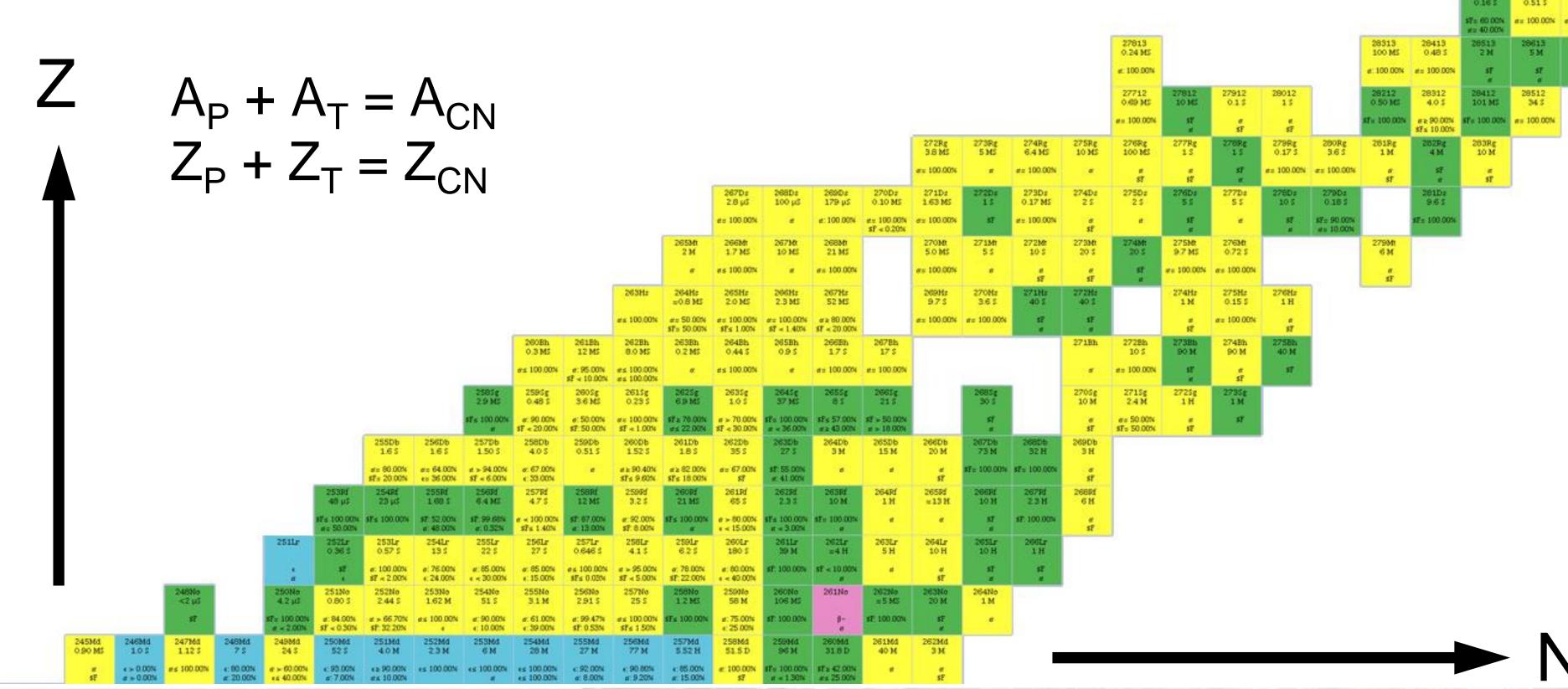
138 nuclides (SHN)



SHN are synthesized in fusion reactions or identified in decay chains of the fusion products

Z = 118

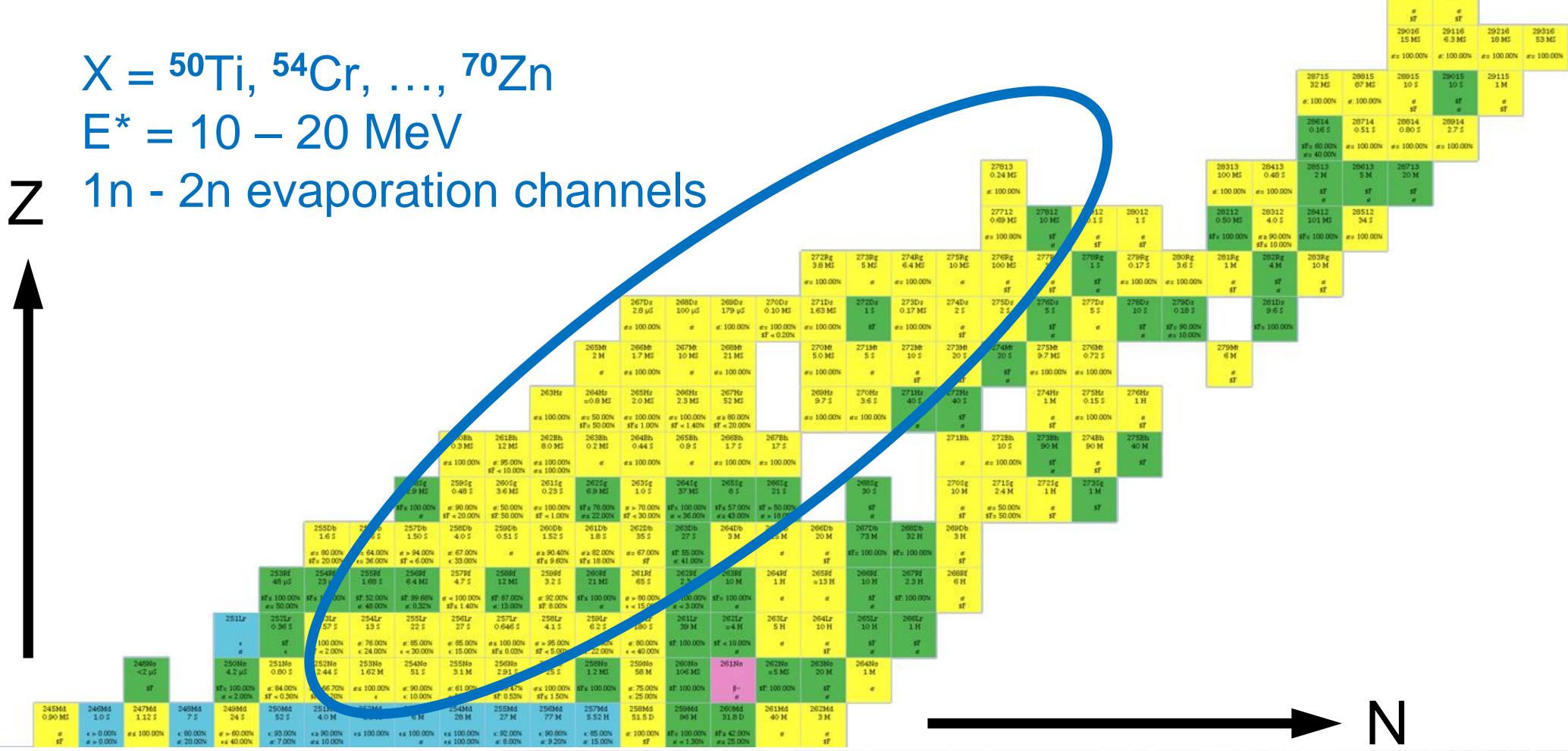
Projectile + Target → SHN*





Cold Fusion Reactions

 $X + {}^{208}Pb, {}^{209}Bi$



Z = 118

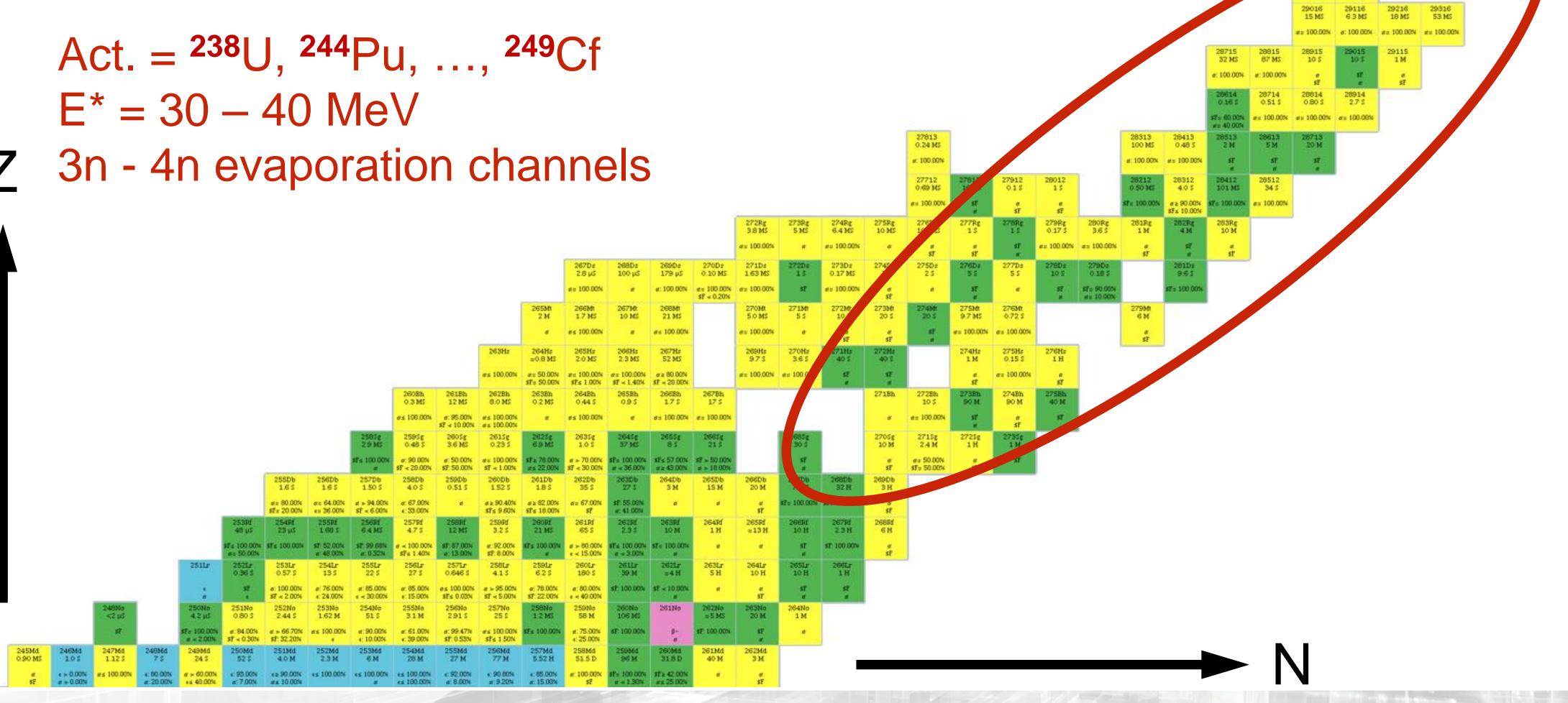
29117 10 MS



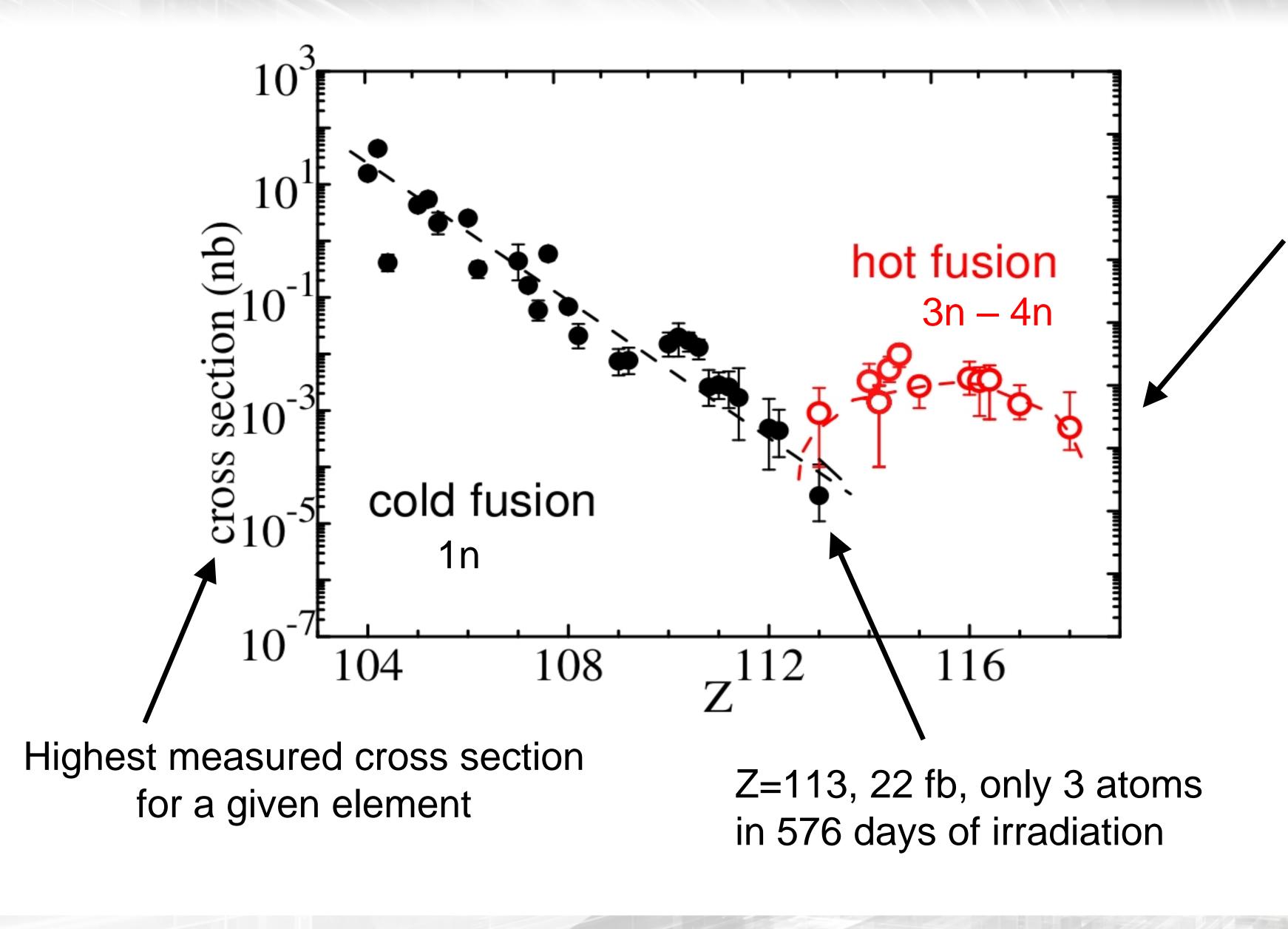
Hot Fusion Reactions

⁴⁸Ca + Actinide targets

Z 3n - 4n evaporation channels







No heavier target than Cf (Z=98) is available

Es (Z=99) is too radioactive

Experiments with ⁵¹V, ⁵⁰Ti, ⁵⁴Cr, and ⁵⁶Fe beams gave no results so far

Most of the experiments are focused on studying SHN properties and (recently) fusion reaction mechanisms.



P_{surv} σ_{ER} **SURVIVAL** COMPOUND P_{fus} **NUCLEUS FUSION** σ_{cap} SYMETRICAL **FISSION CAPTURE FAST FISSION**

FBD (fusion-by-diffusion)

Synthesis of SHN can be described as a 3 step process:

$$\sigma_{ER} = \sigma_{cap} \; P_{fus} \; P_{surv}$$

Not measured directly Difficult to calculate

n be dels

> Smoluchowski Diffusion equation

Well established theory and formulas



masses, fission
barriers, deformations
from Warsaw
Micro-Macro model

W. J. Świątecki, K. Siwek-Wilczyńska,

J. Wilczyński, PRC 2005

T. Cap et al., PRC 2011

K. Siwek-Wilczyńska et al. PRC 2012

T. Cap et al., PRC 2013

K. Siwek-Wilczyńska et al. PRC 2019

Measured in experiments, can be calculated using various models



Diffussed barrier formula (Gaussian distribution around mean entrance channel barrier *B0*)

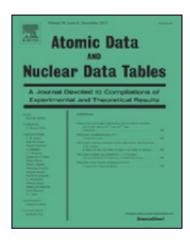




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Properties of heaviest nuclei with $98 \le Z \le 126$ and $134 \le N \le 192$



P. Jachimowicz a, M. Kowal b,*, J. Skalski b

Ground-state and saddle-point shapes and masses for 1305 heavy and superheavy nuclei including odd-A and odd–odd systems. Static fission barrier heights, one- and two-nucleon separation energies, and Qα values.

Microscopic–macroscopic method with the deformed Woods–Saxon single-particle potential and the Yukawa-plus-exponential macroscopic energy taken as the smooth part.

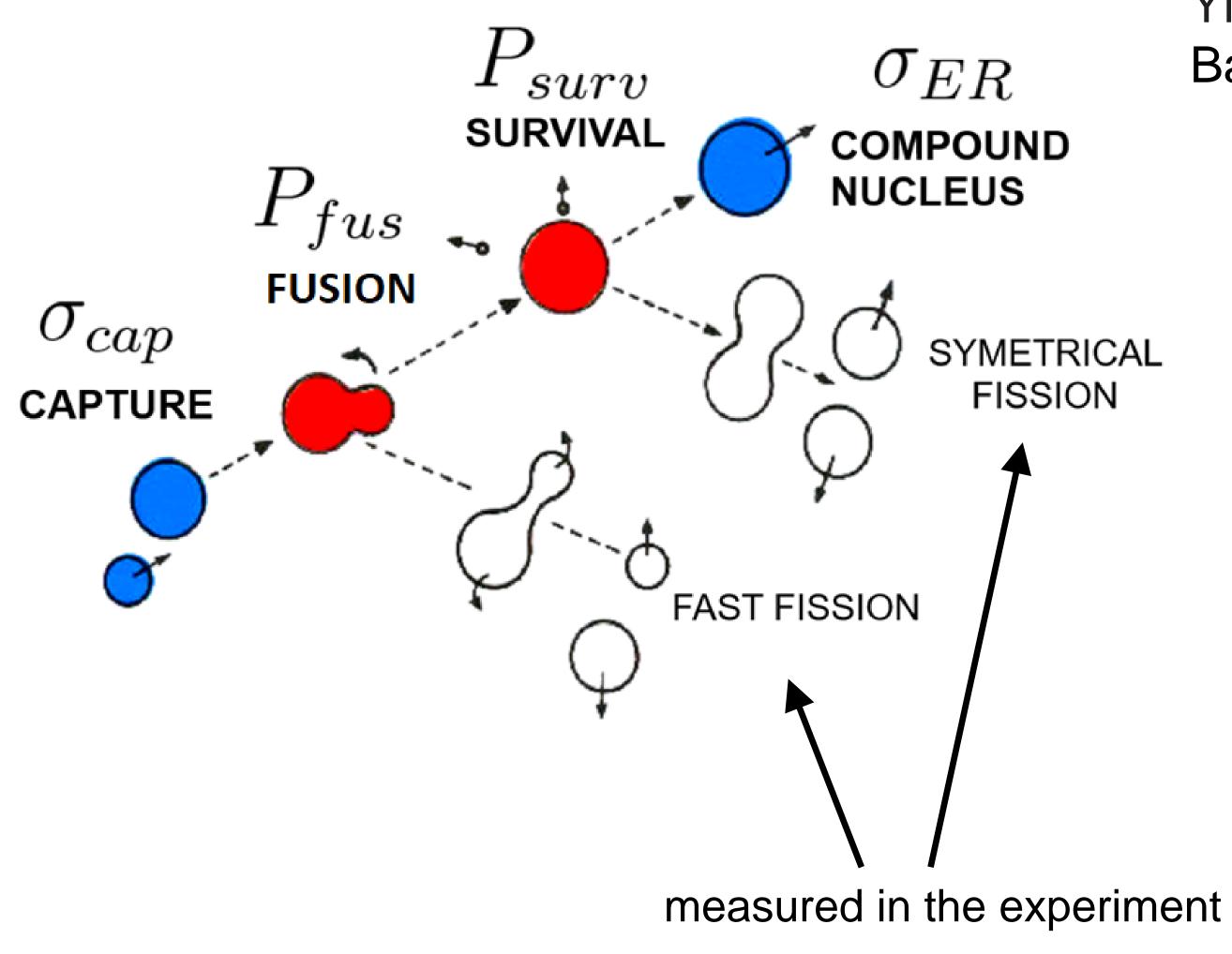
Ground-state shapes and energies are found by the minimization over **seven axially-symmetric deformations**. A search for saddle-points was performed by using the "imaginary water flow" method in three consecutive stages, using five- (for nonaxial shapes) and seven-dimensional (for reflection-asymmetric shapes) deformation spaces.

Good agreement with the experimental data for actinides.

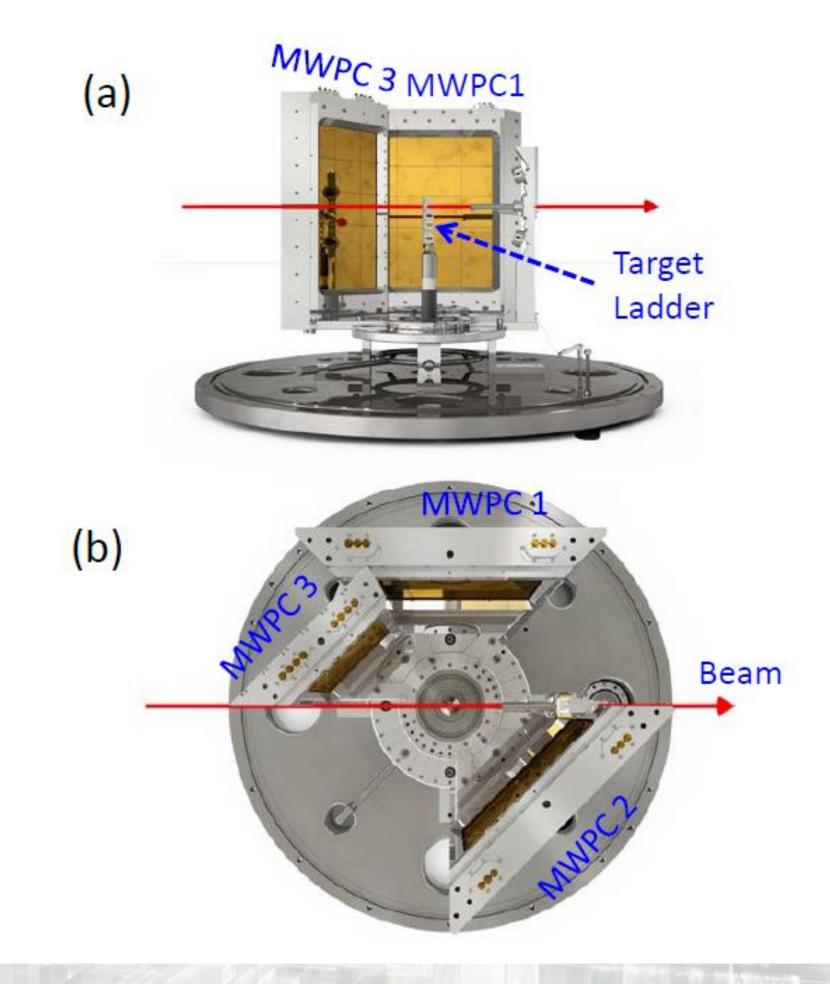


^a Institute of Physics, University of Zielona Góra, Szafrana 4a, 65-516 Zielona Góra, Poland

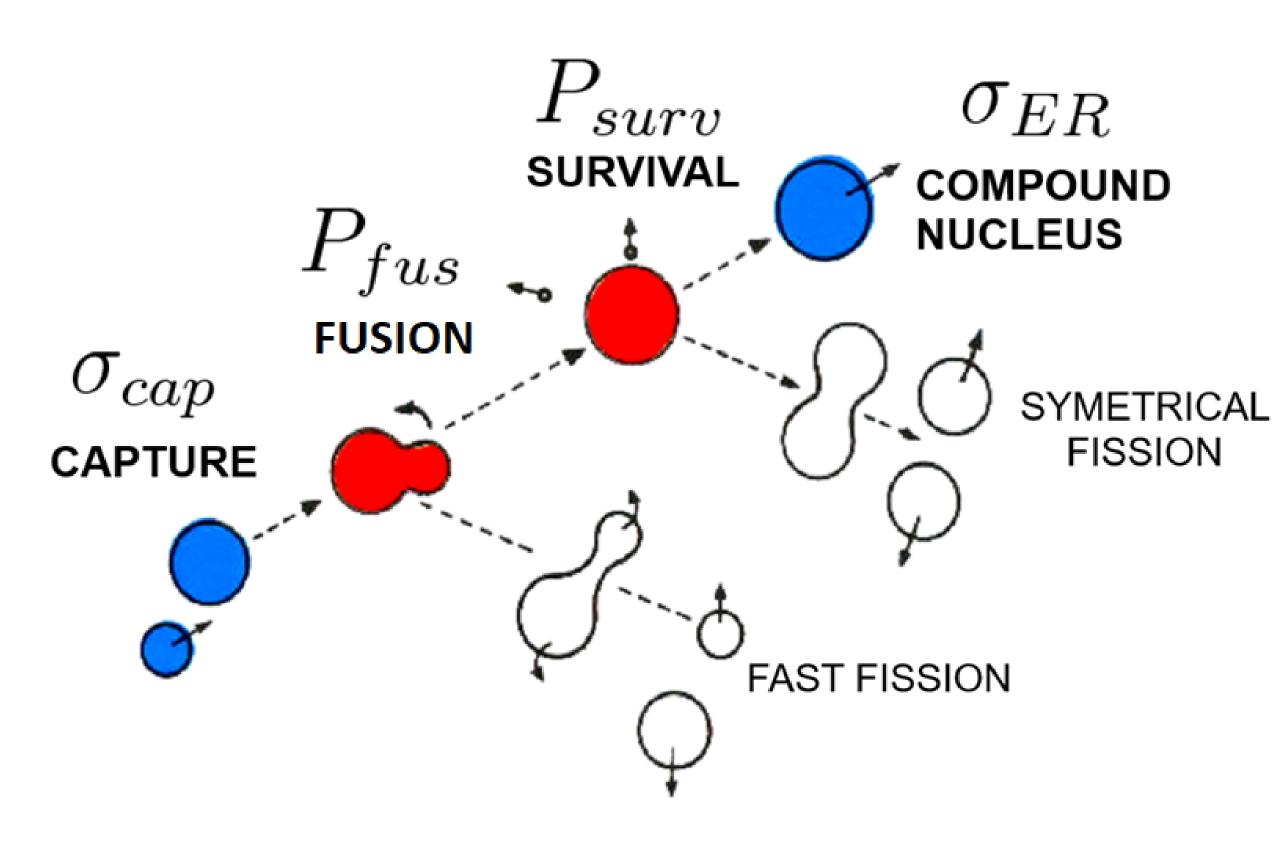
^b National Centre for Nuclear Research, Pasteura 7, 02-093 Warsaw, Poland



Mechanisms Suppressing Superheavy Element Yields in Cold Fusion Reactions Banerjee et al., PRL 122, 232503 (2019)



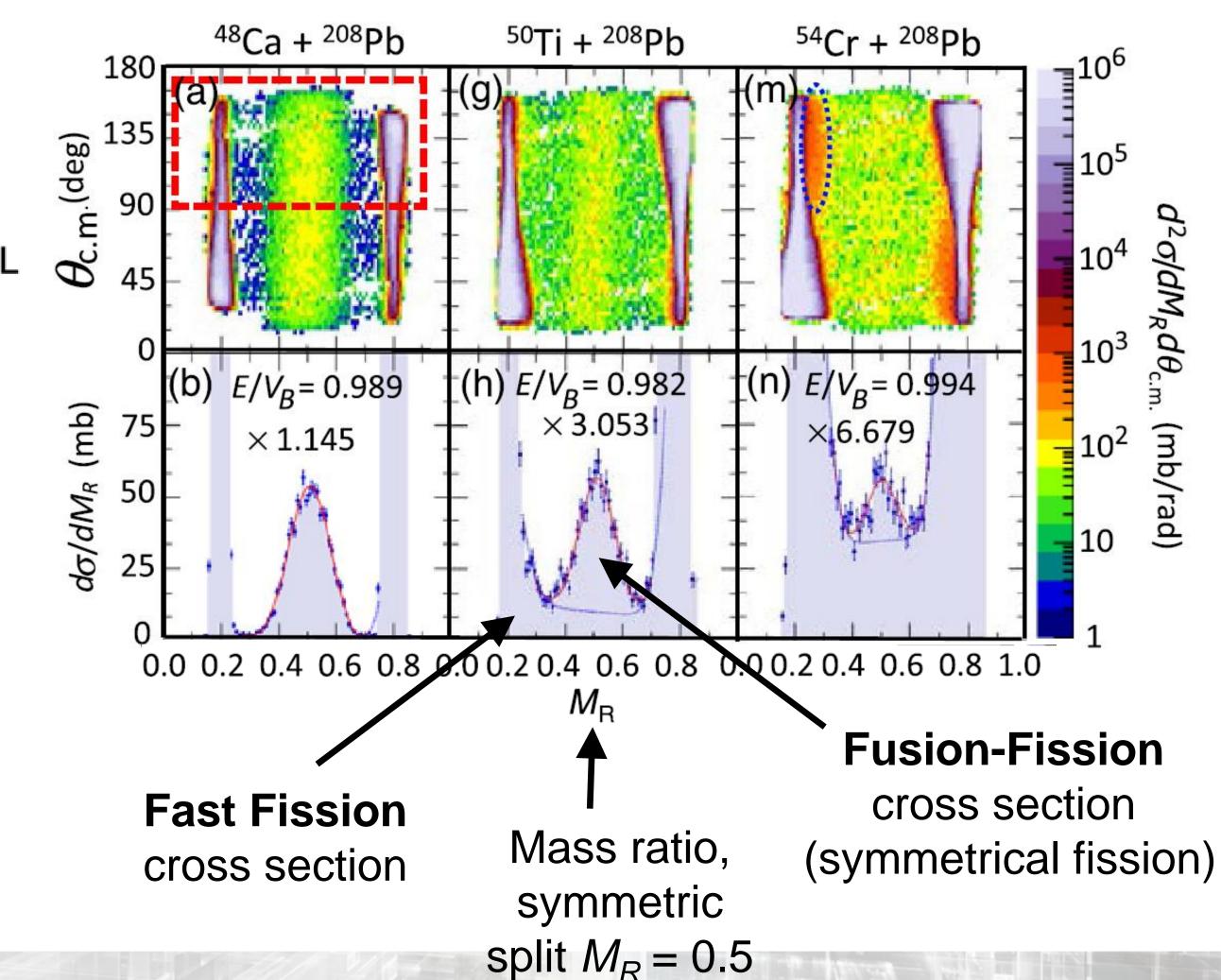




Pfus can be experimentally estimated:

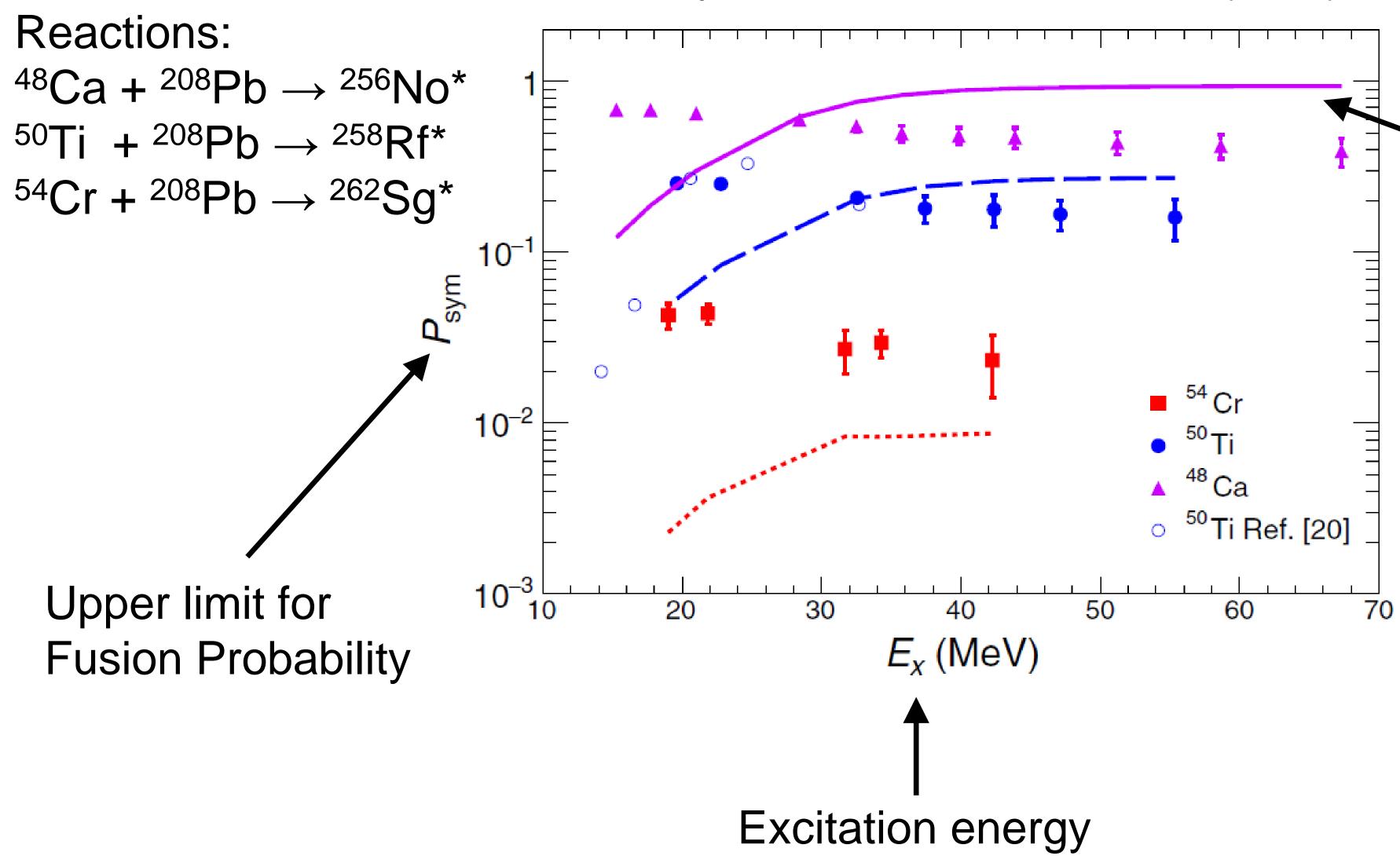
$$Psym = \frac{\textbf{Fusion-Fission cross section}}{\textbf{Capture cross section}}$$

Mechanisms Suppressing Superheavy Element Yields in Cold Fusion Reactions Banerjee et al., PRL 122, 232503 (2019)





Banerjee et al., PRL 122, 232503 (2019)



Difusion model calculations by V. Zagrebaev and W. Greiner PRC 78, 034610 (2008).

The experimental trend is different than the model predictions for all 3 reactions.

The conclusion was that diffusion is not the main mechanism responsible for the synthesis of SHN.

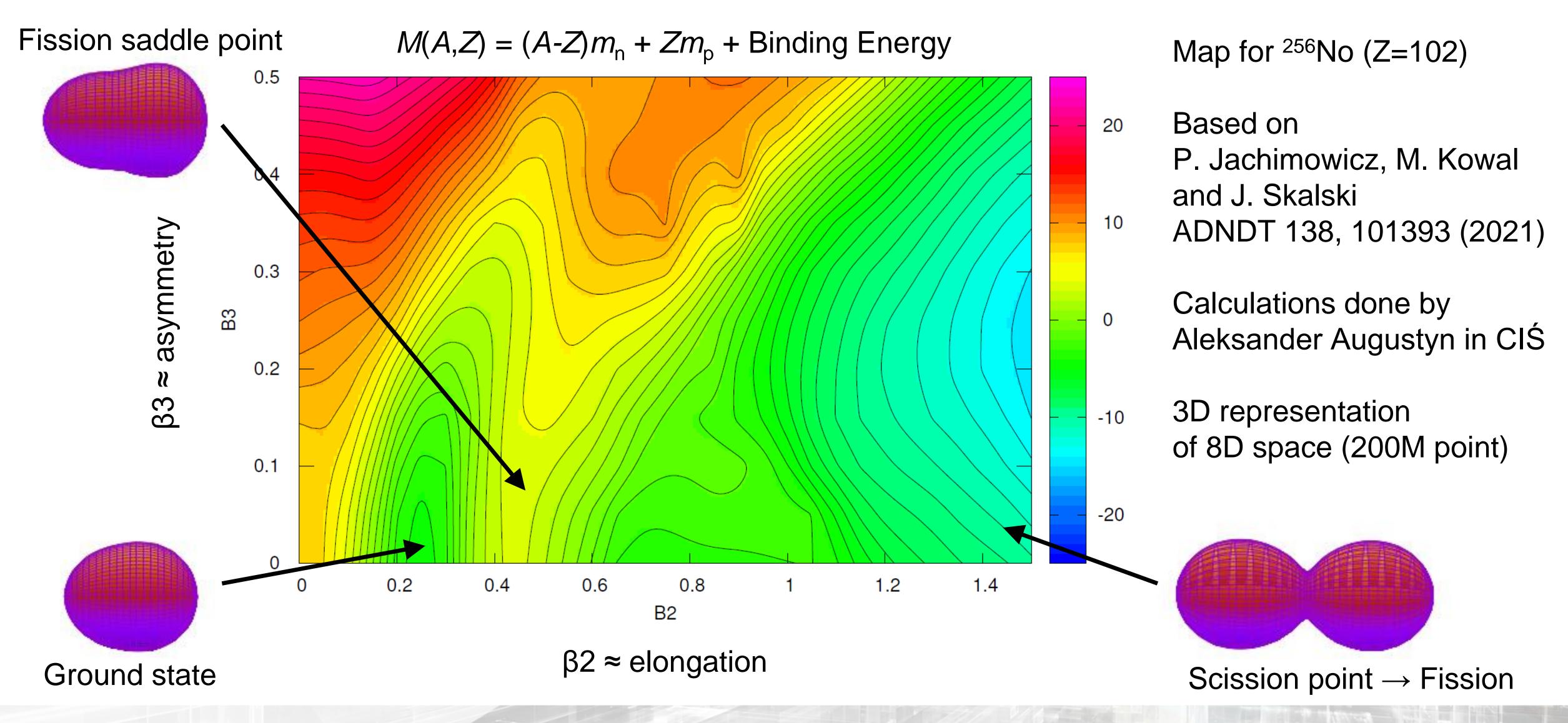
Questions:

Can the diffusion approach (FBD) describe the experimental results?

What is the fusion mechanism?

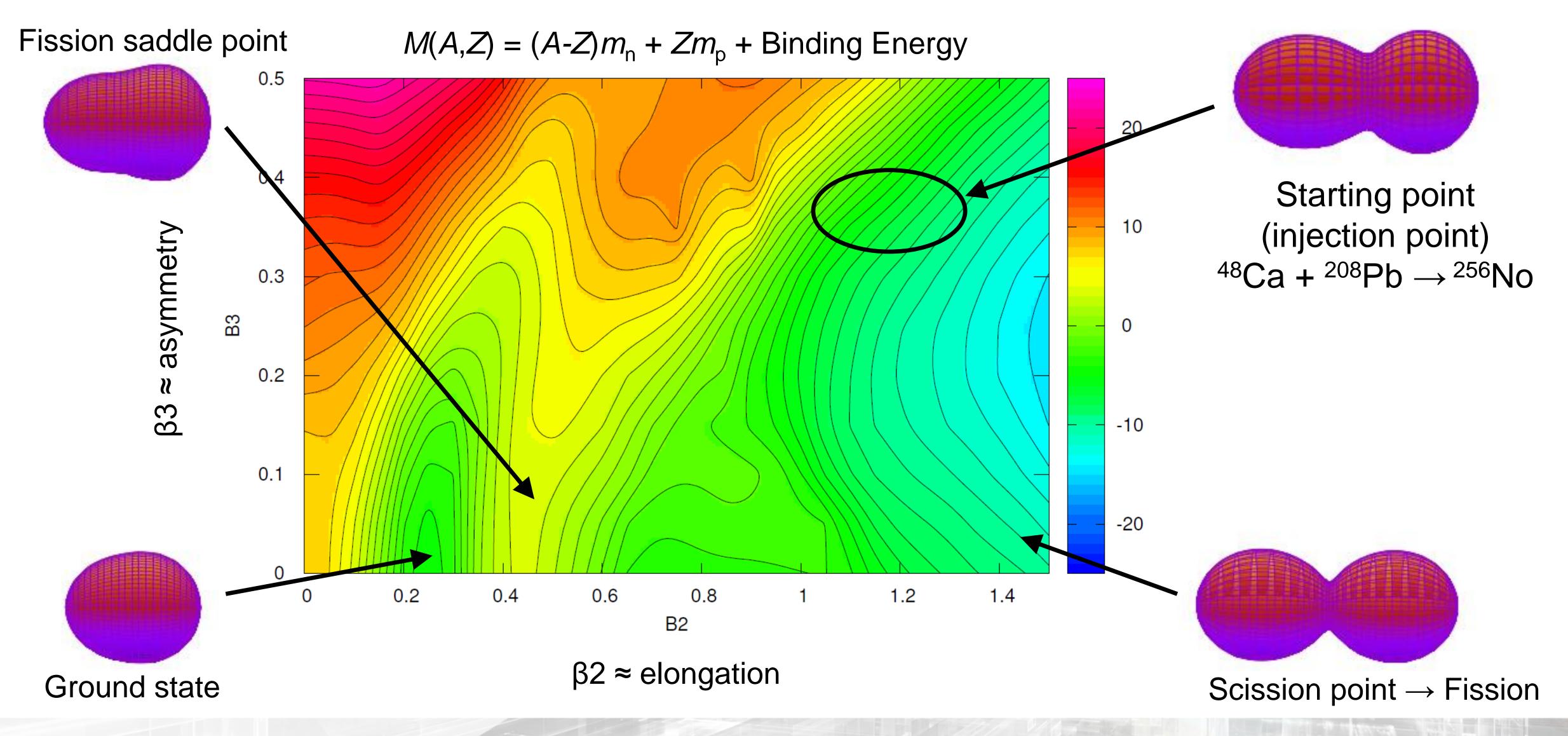


Fusion process on the Potential Energy Surface (PES)



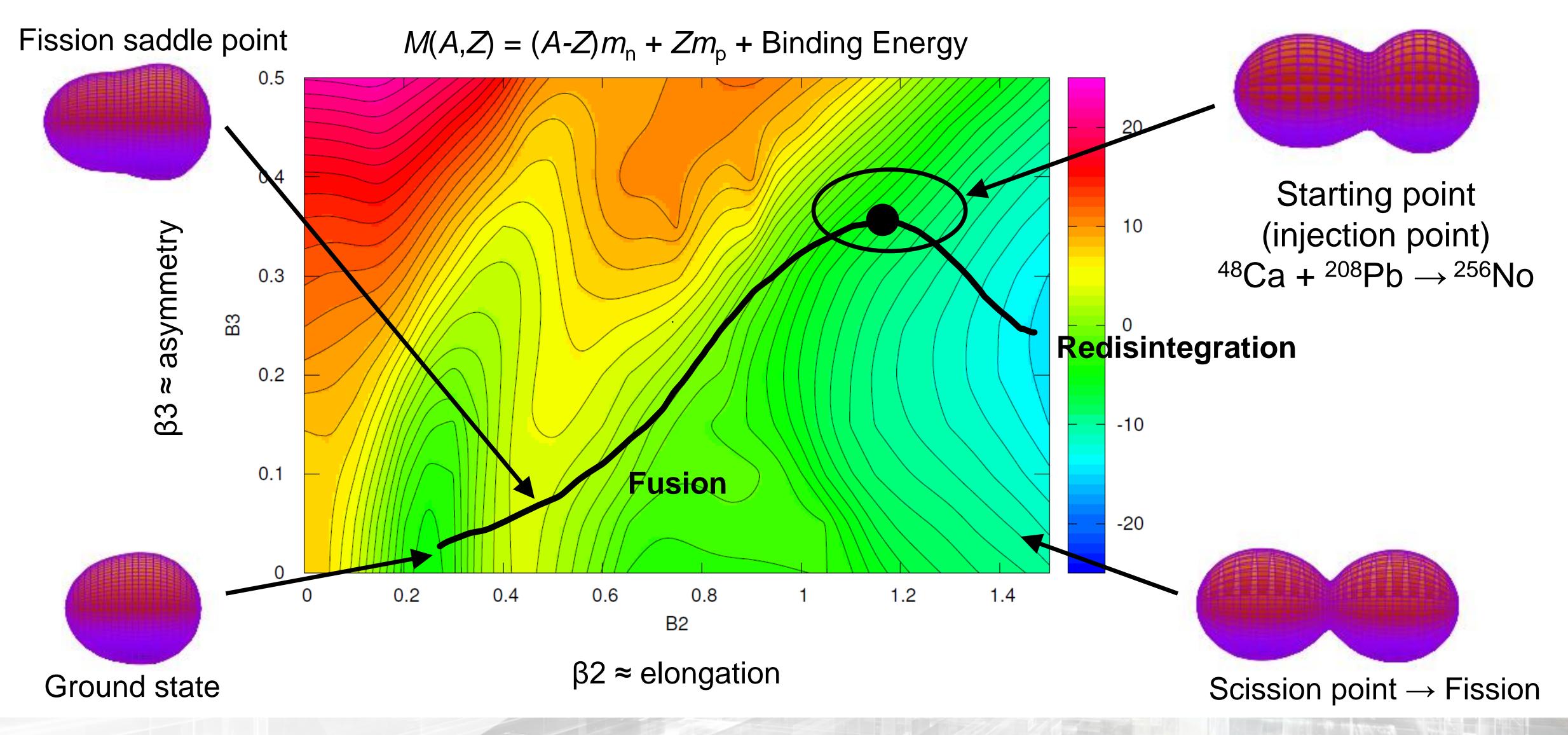


Fusion process on the Potential Energy Surface (PES)





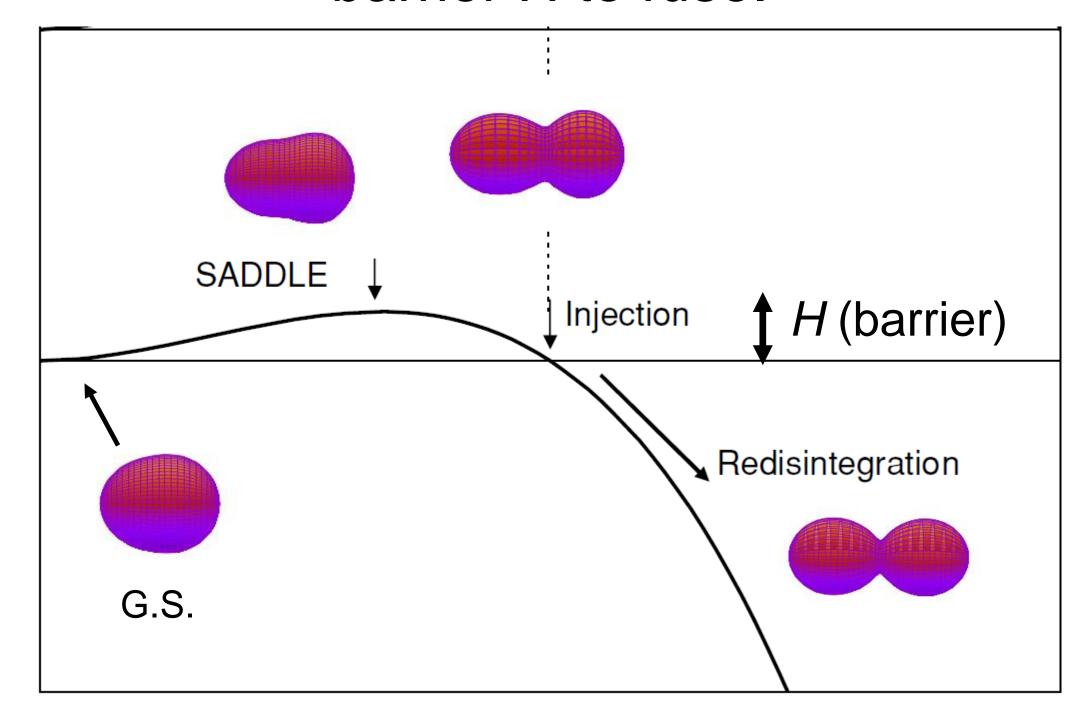
Fusion process on the Potential Energy Surface (PES)





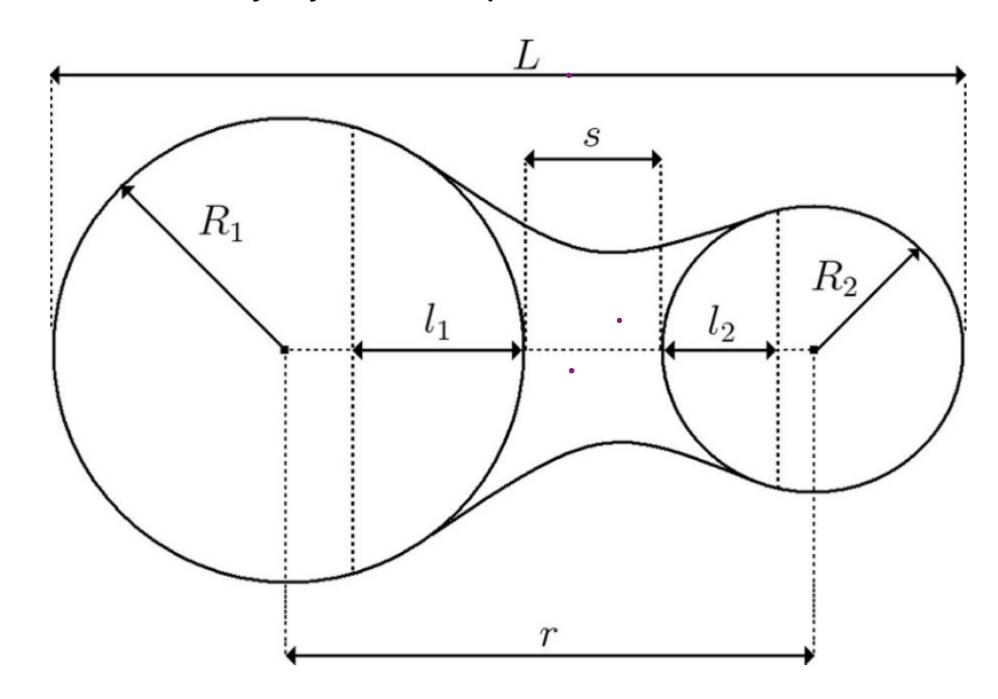
In the FBD model we use 1D motion approximation

The system must overcome an internal barrier *H* to fuse.



L is the effective elongation (along the fusion path)

Macroscopic deformation energies are calculated using the parameterization of the nuclear shapes by two spheres joined smoothly by a third quadratic surface of revolution.



The distance between the nuclear surfaces of two colliding nuclei at the **injection** point, s_{inj} , is the only adjustable parameter of the model.

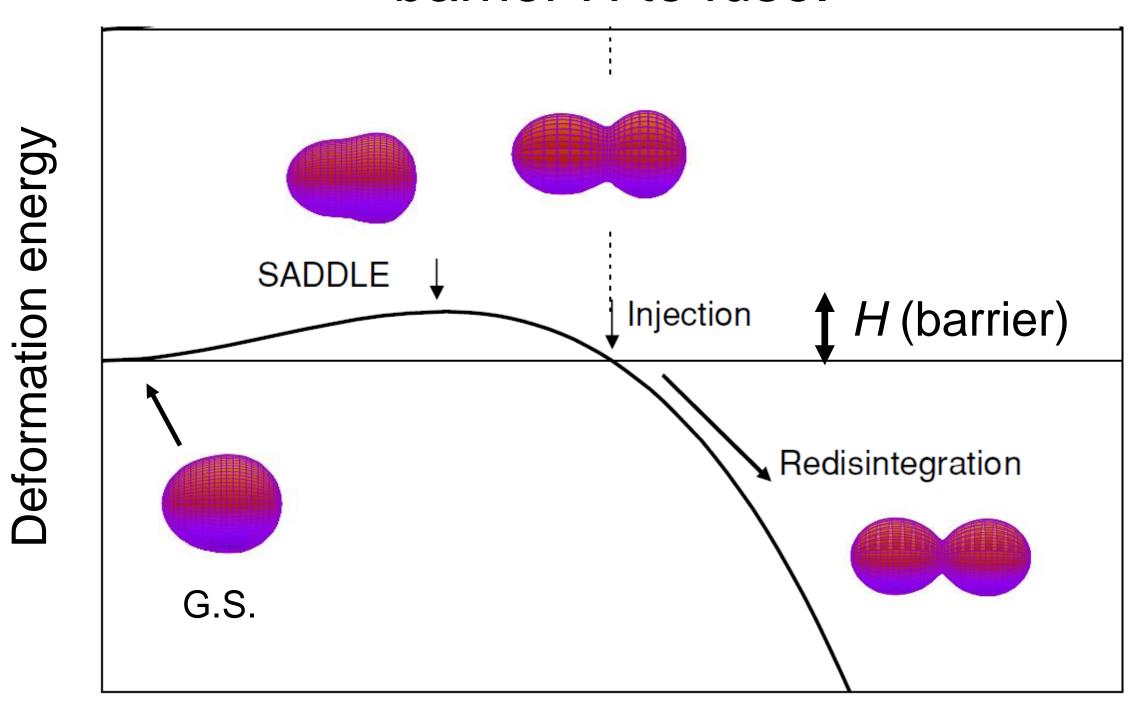
 s_{inj} distance was parametrized by analyzing 27 cold fusion reactions.



Deformation energy

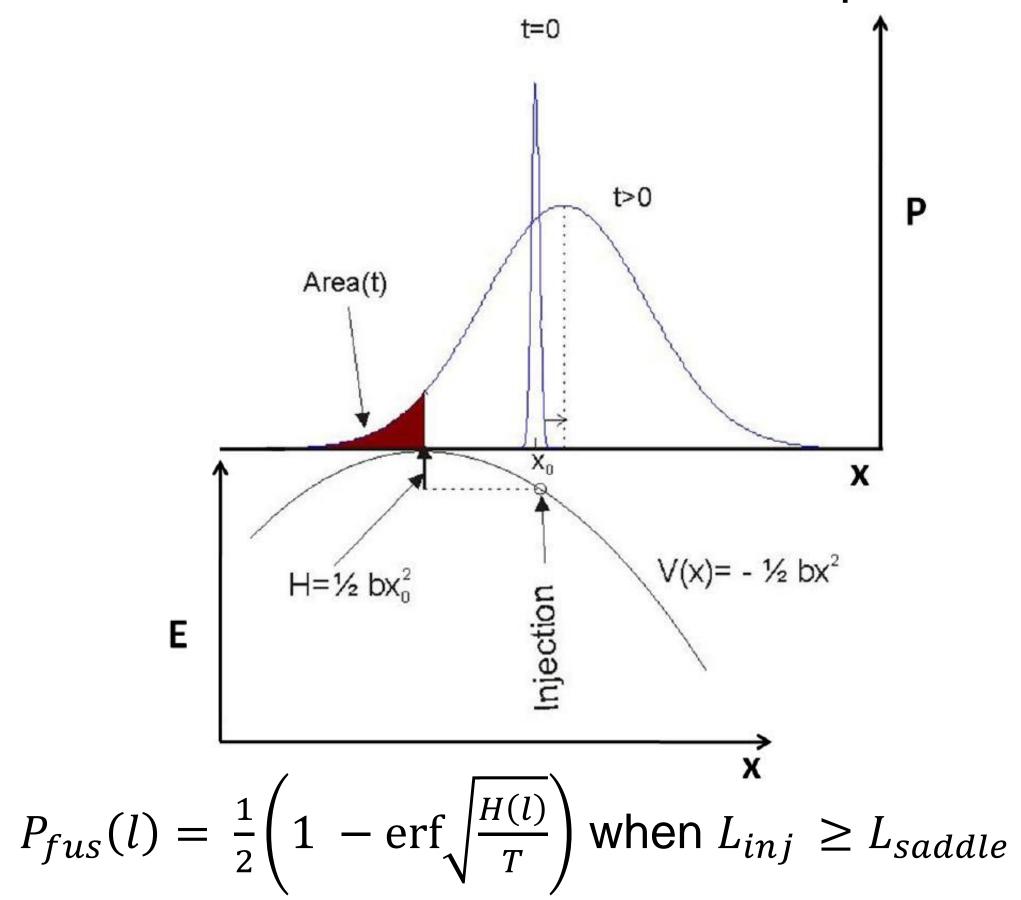
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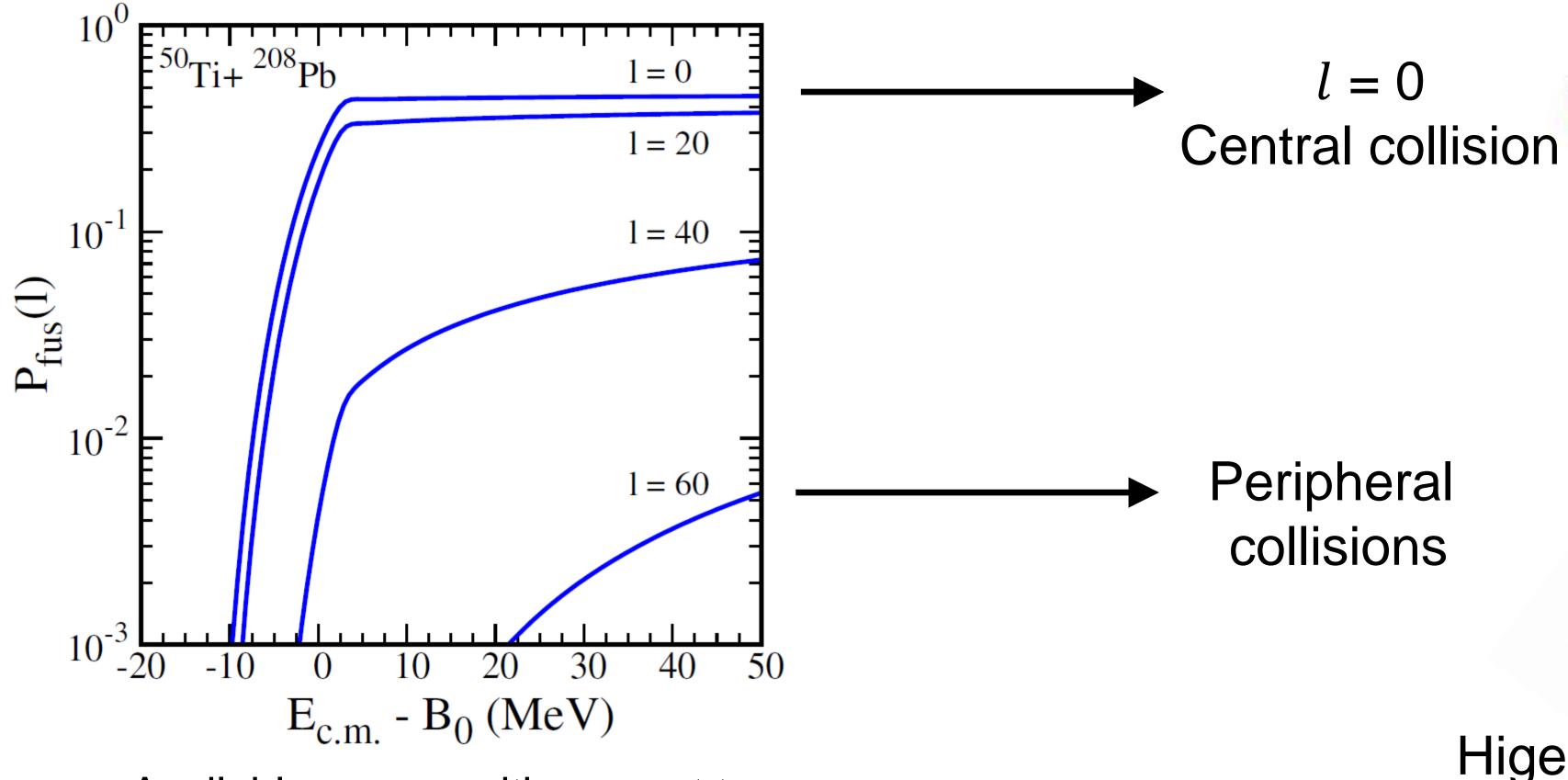
Pfus is calculated by solving 1D Smoluchowski Diffusion Equation



H(l) – the function of angular momentum and bombarding energy

T – the temperature depends on available energy

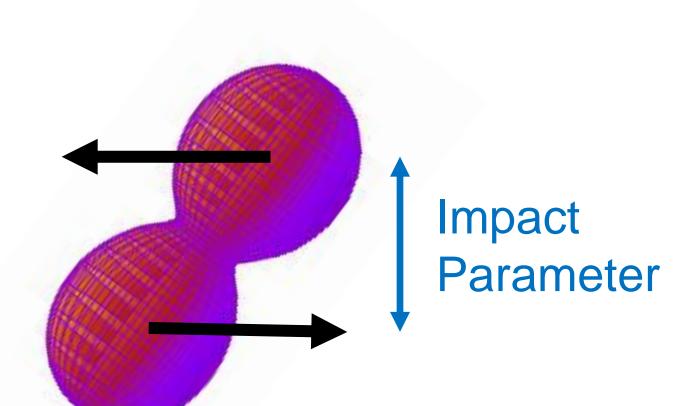






$$B0 = 191.3 \text{ MeV}$$

 $E^*(E_{c.m.} = B0) = 21.7 \text{ MeV}$



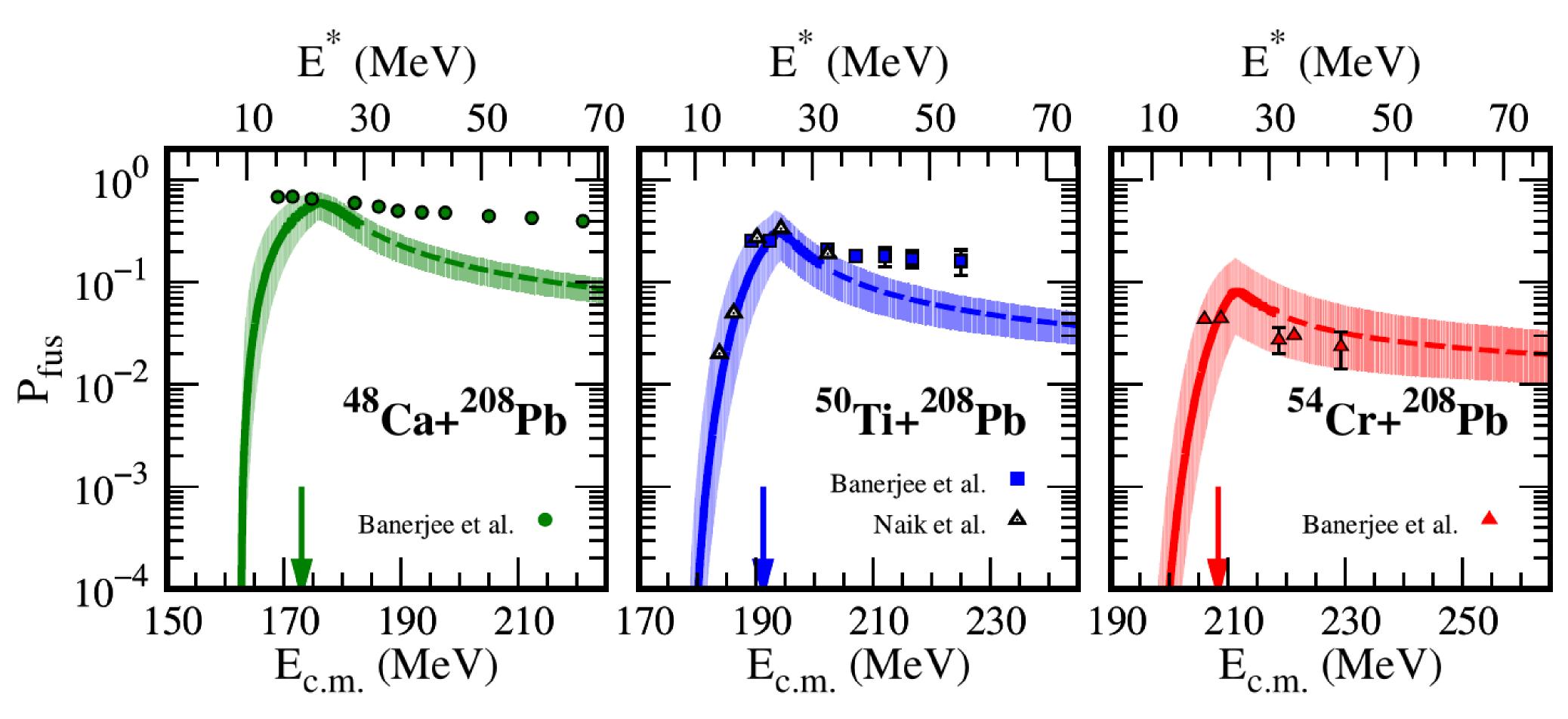
Higer partial waves *l* =

Higher rotational Energy

Lower Pfus



$$P_{fus}(E_{c.m.}) = \frac{\sum_{l=0}^{lmax} (2l+1)P_{fus}(l)}{(2l_{max}+1)^2}$$
 Fusion probability averaged over l



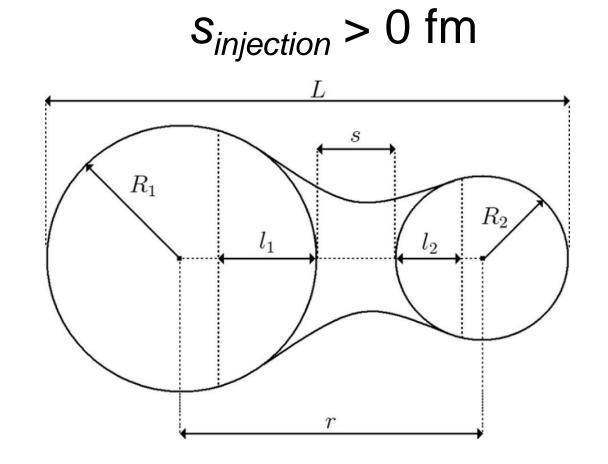
Submitted to PLB, available at: arXiv:2107.00579

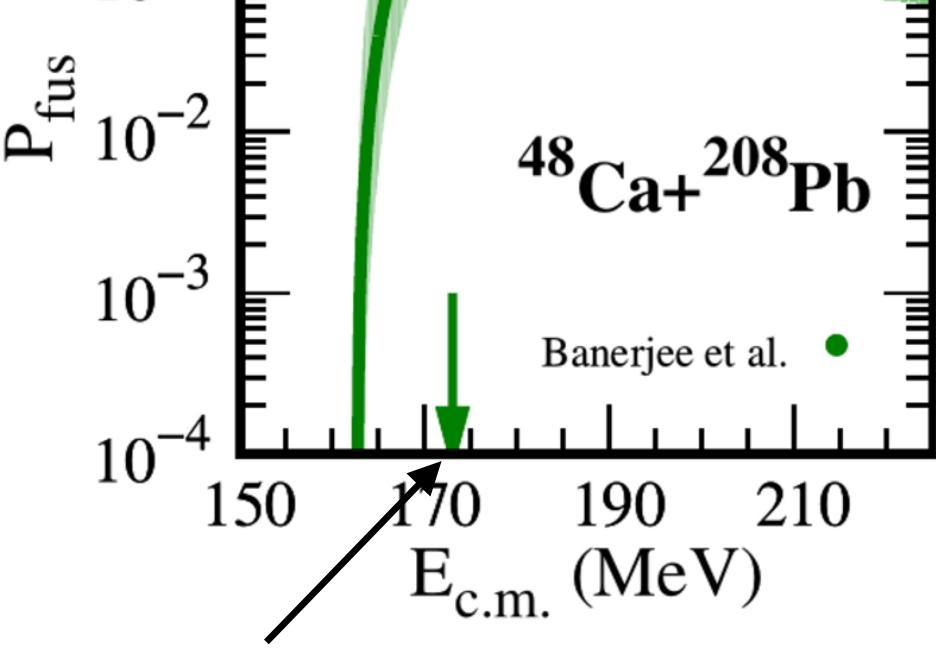


Tangent configuration of projectile and target $(s_{injection} = 0 \text{ fm})$

Below *B0*, the *Pfus* growth comes from the reduction in the height of the internal barrier opposing fusion.

The *Pfus* saturation above *B0* results from suppression of the contributions from higher partial waves and can be linked to the critical angular momentum.



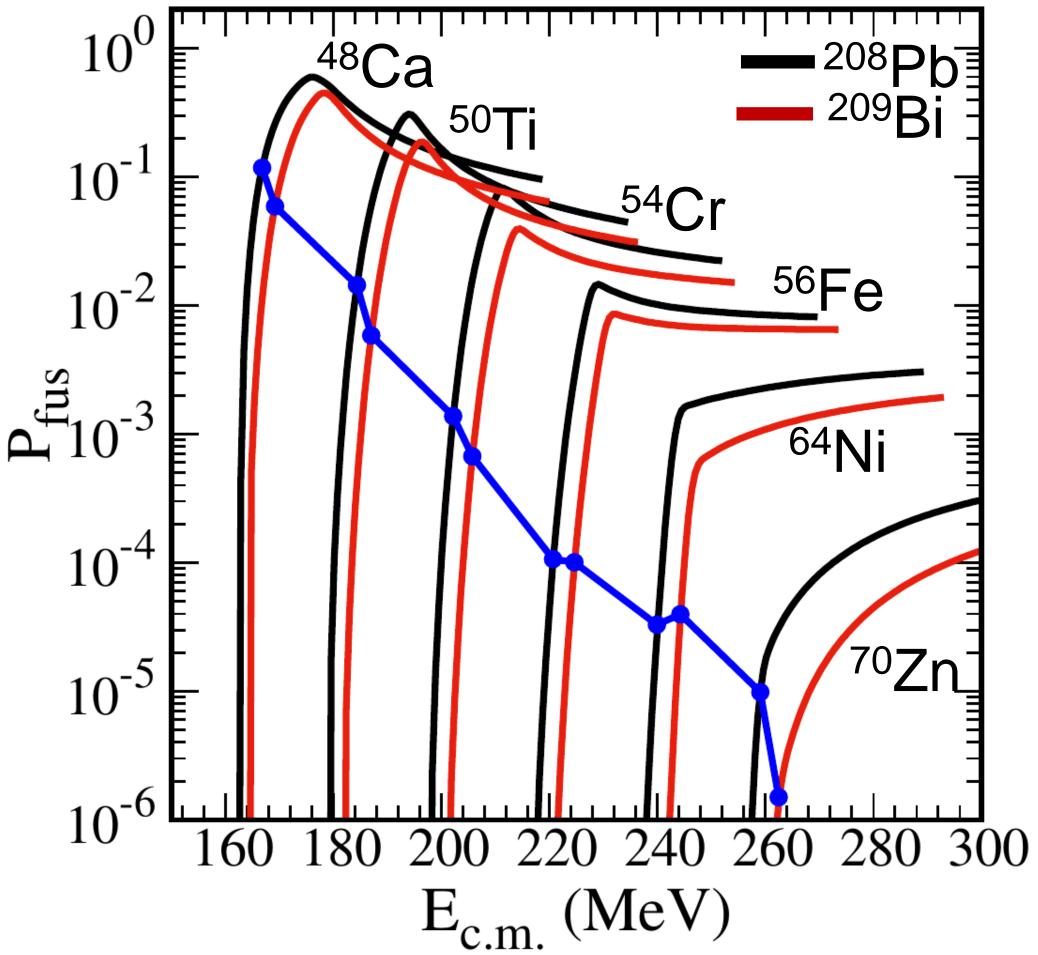


E* (MeV

The difference between rotational energies in the fusion saddle and the contact (sticking) configuration plays a major role in compound nucleus formation at energies above *B0*.

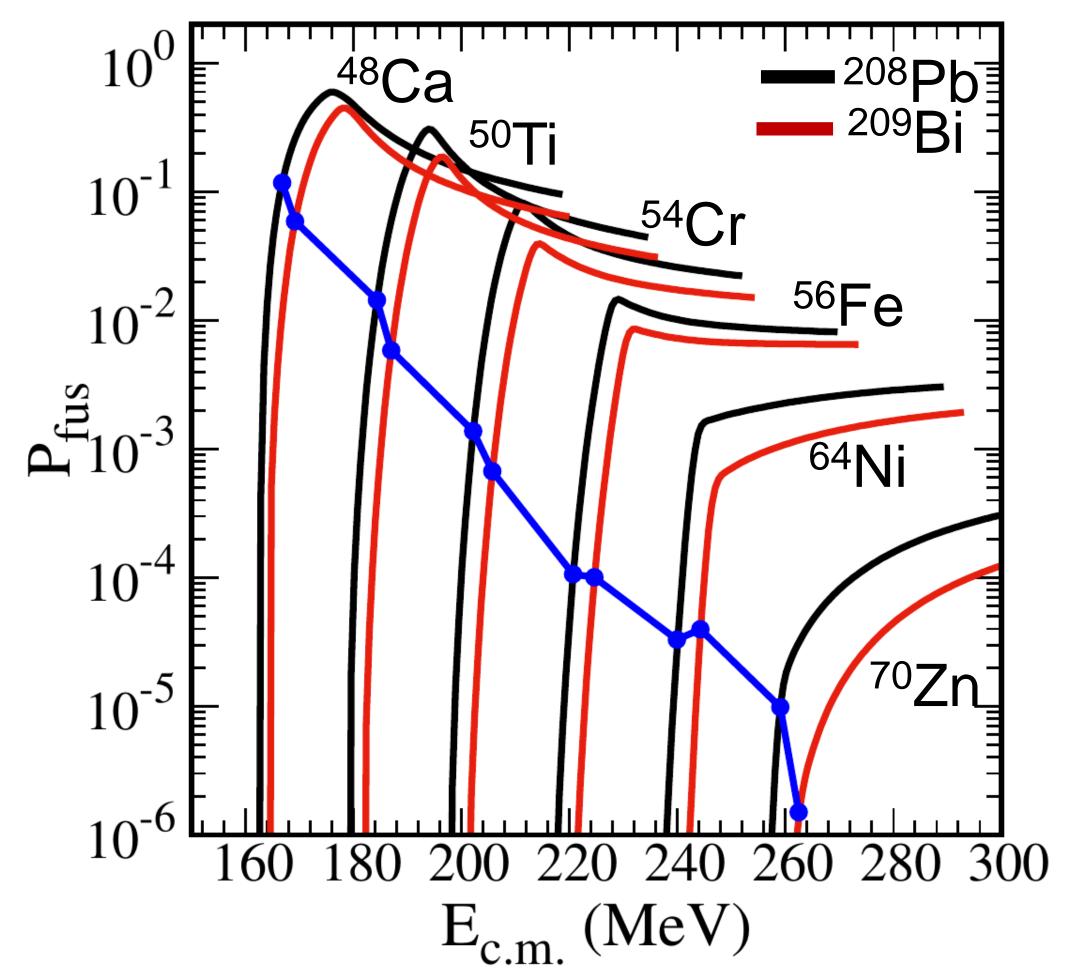
B0 - entrance channel barrier (Coulomb+Nuclear potential)



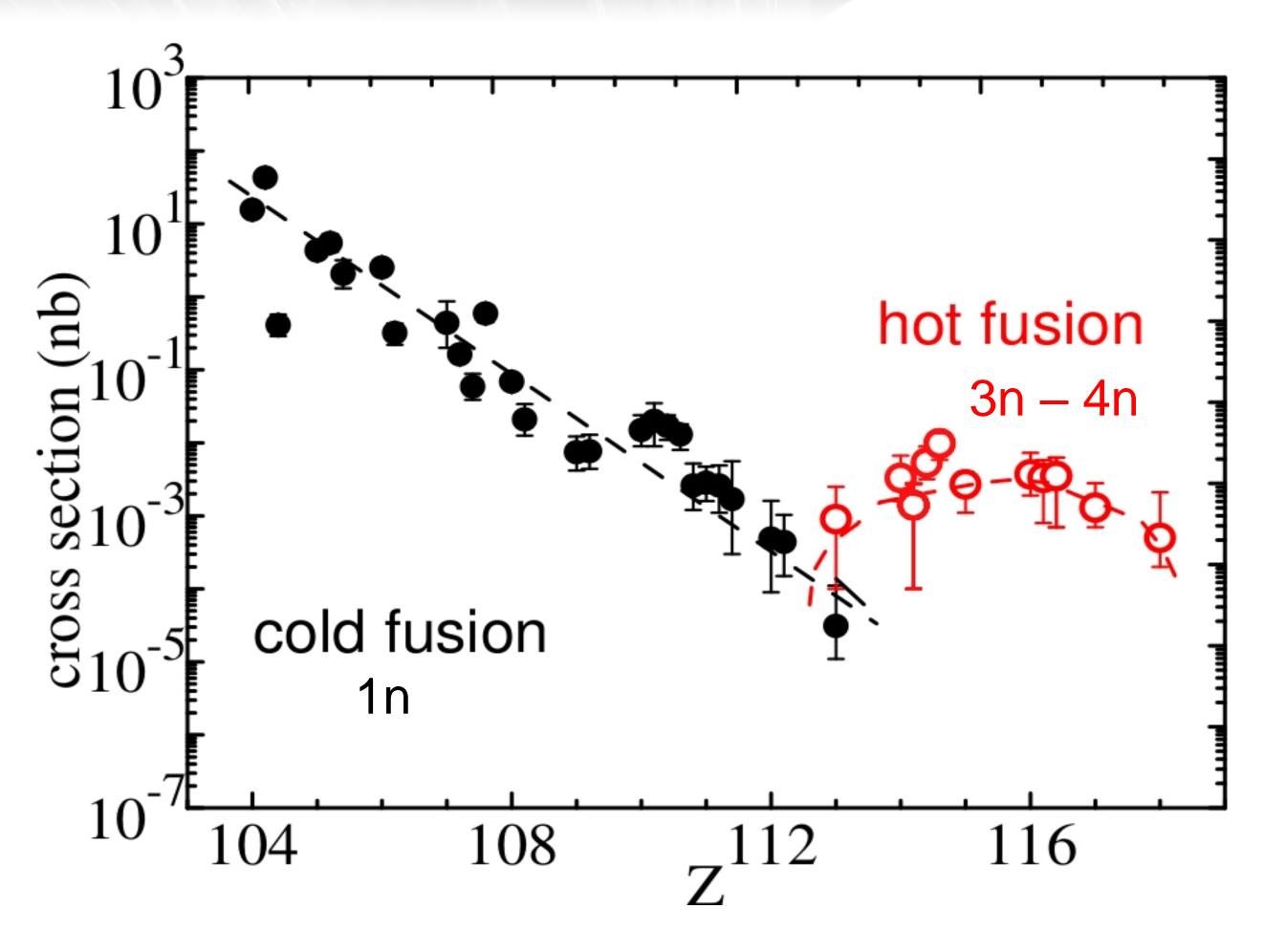


Blue line – Maxima of the 1n channel cross section for cold Fusion reactions



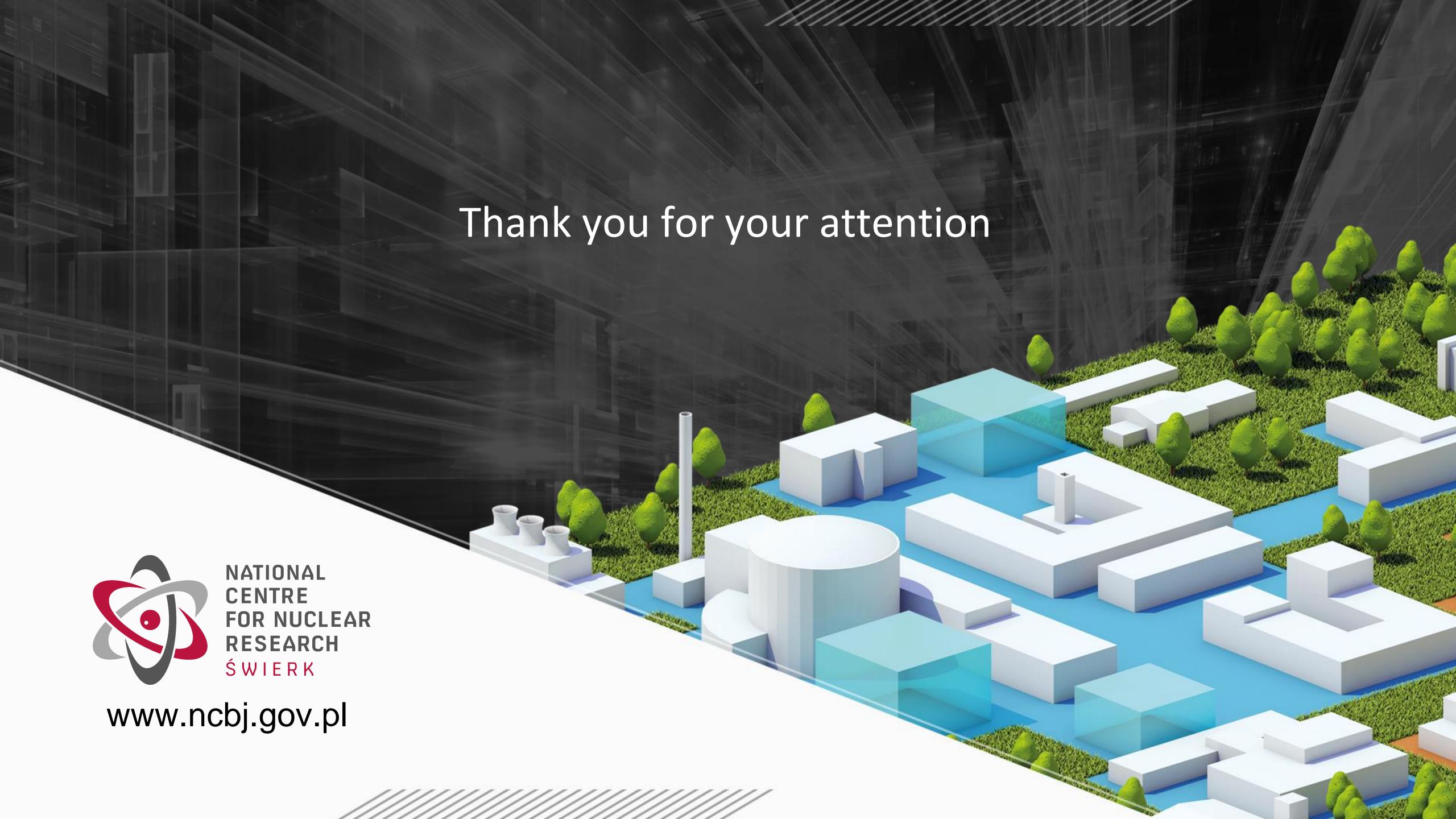


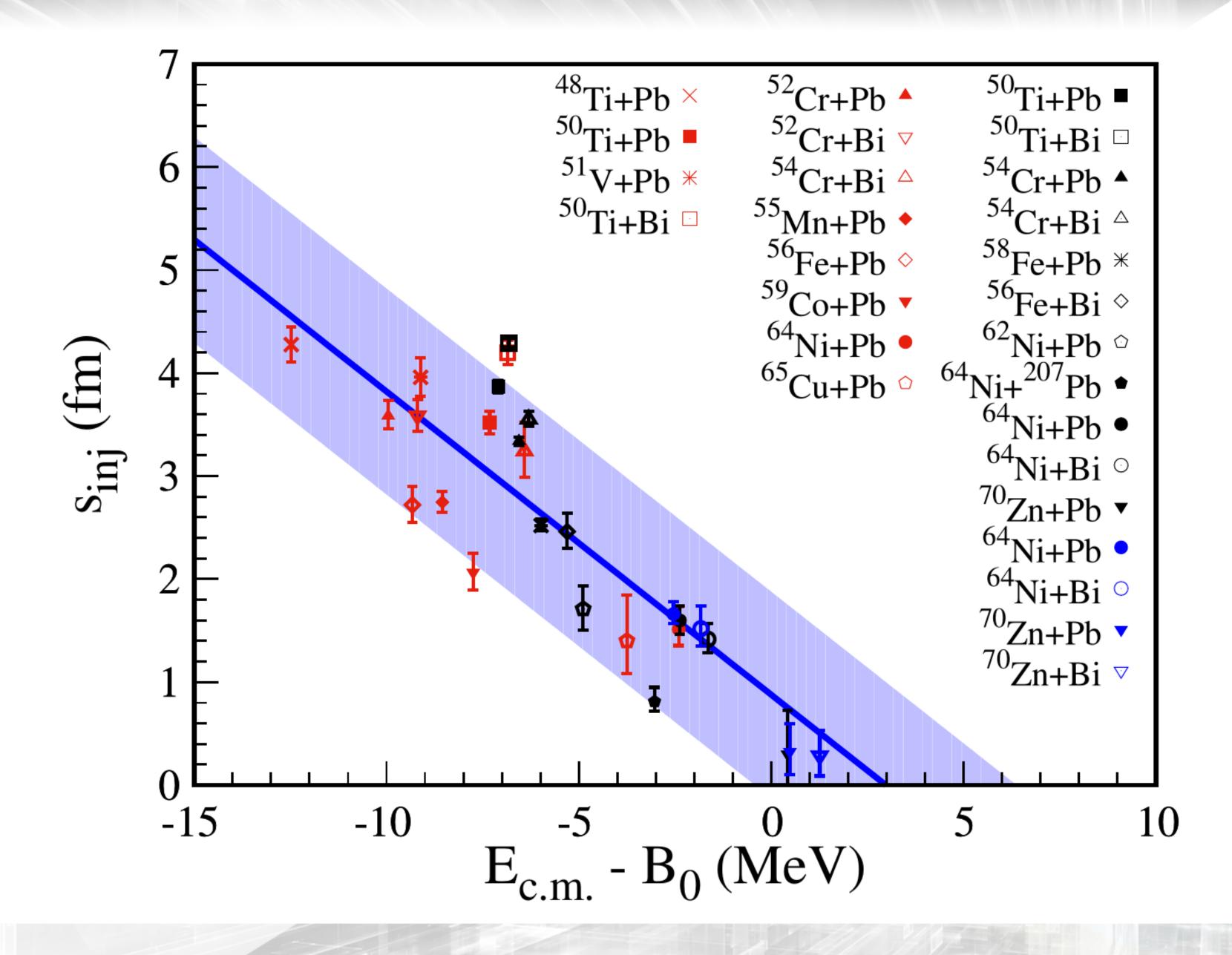
Blue line – Maxima of the 1n channel cross section for cold Fusion reactions



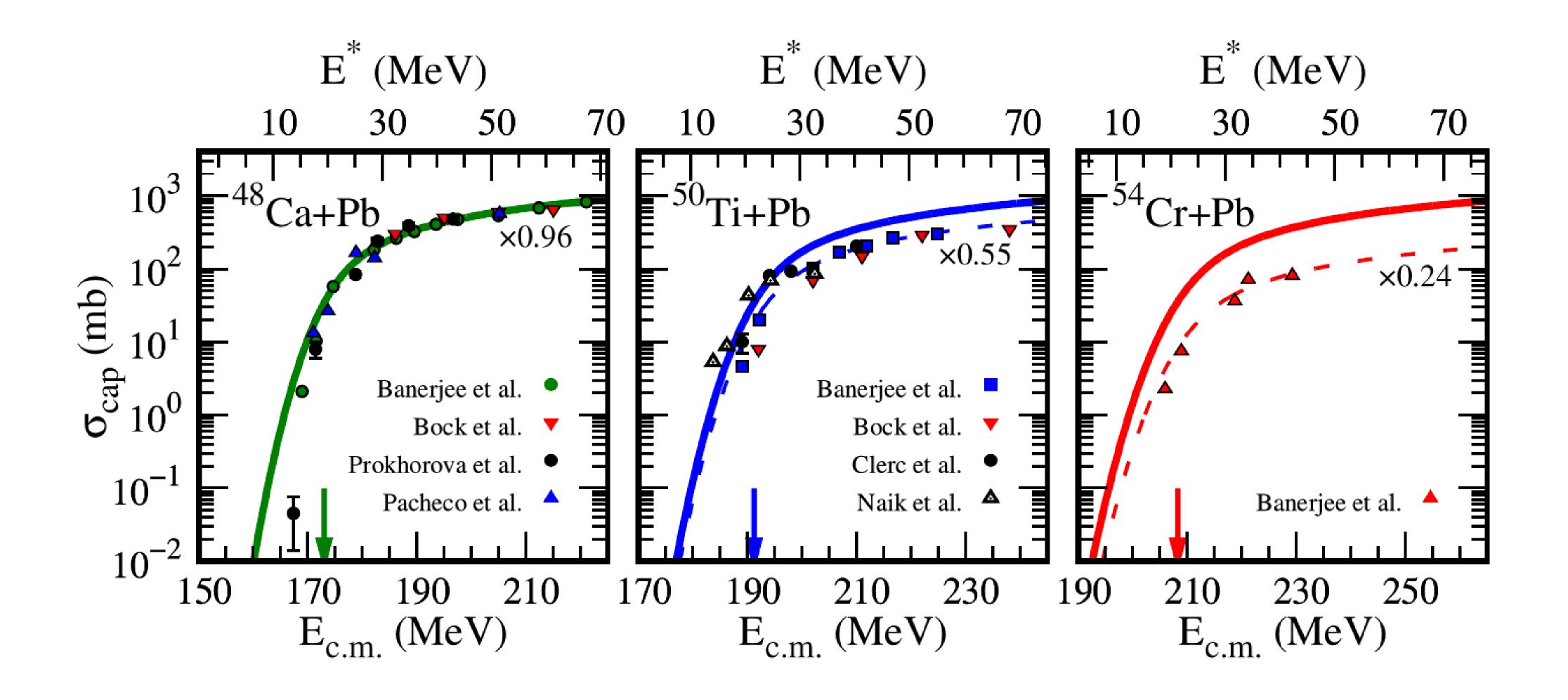
Next step – calculations for hot fusion reactions











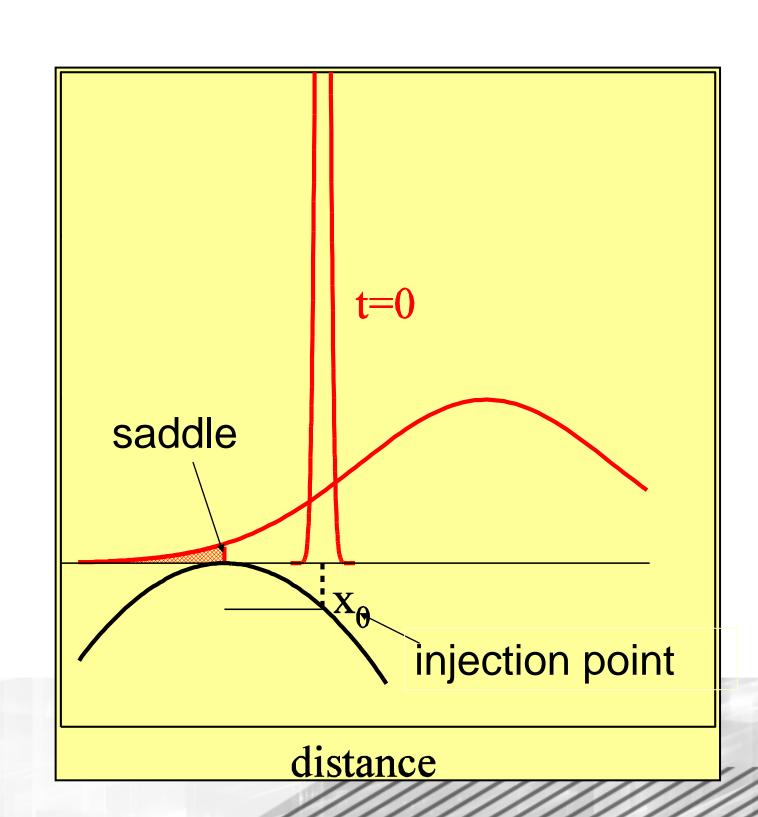


Smoluchowski Diffusion Equation

$$\gamma \frac{\partial W(x,t)}{\partial t} = -\frac{\partial}{\partial x} \left[\underline{bx} W(x,t) \right] + T \frac{\partial^2}{\partial x^2} W(x,t)$$

$$\propto \text{ Viscosity of fluid } \text{ Driving force } \text{ Temperature}$$

W(x,t) = probability to find Brownian particle at position x at time t



Exact solution in a parabolic potential $V(x) = -bx^2/2$ is a sliding, swelling Gaussian.

P(fusion) = fraction of Gaussian captured inside the barrier as $t\rightarrow\infty$

P(fusion) =
$$\frac{1}{2}(1-\text{erf}\sqrt{B/T})$$
 B = $bx_0^2/2$ if $x_0 \ge 0$ (injection point),

W.J. Świątecki, K. Siwek-Wilczyńska, J. Wilczyński Acta Phys. Pol. B34 (2003)2049, IJMP E13 (2004) 261, Phys.Rev.C71 (2005) 014602





